

Efficient G_0W_0 using localized basis sets: a benchmark for molecules

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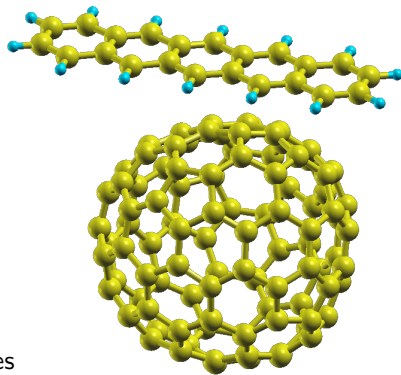
NanoSpain conference
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<http://www.nanospainconf.org/2017>

Prerequisite for computer-aided nano structures design

- ! *LARGE* [number of atoms] usually implies *MANY* calculations, i.e. the calculations are better to be *FAST*.
- ▶ LCAO provides a basis for ab-initio, economical description of electronic structure; *ANALYSIS* is mostly straightforward.
- ▶ Locality is built-in into LCAO explicitly, allowing for *FAST* algorithms to be developed.
- ▶ This talk is about our efforts towards a fast, quantitative GWA

Outline of the talk

- ▶ Theory of GWA : simplistic
- ▶ Methods to realize GWA
- ▶ IPs for G2/97 test set
- ▶ IPs & EAs for acceptor molecules
- ▶ Analysis of core electron contributions to G_0W_0
- ▶ Conclusions



GWA as an extension of Hartree-Fock method

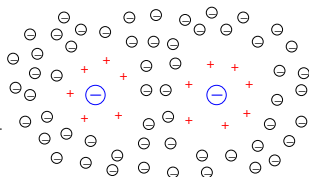
$$\text{Fock exchange: } \Sigma_x(\mathbf{r}, \mathbf{r}') = - \sum_{n \in \text{occ}} \frac{\psi_n(\mathbf{r})\psi_n(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|}$$

GW self-energy

$$\Sigma(\mathbf{r}, \mathbf{r}', \omega) = \frac{i}{2\pi} \int_{-\infty}^{\infty} d\omega' G(\mathbf{r}, \mathbf{r}', \omega + \omega') W(\mathbf{r}, \mathbf{r}', \omega')$$

$$W = v - v \chi_0 W.$$

$$\chi_0(\mathbf{r}, \mathbf{r}', \omega) = \sum_{n,m} \frac{\psi_n(\mathbf{r})\psi_m(\mathbf{r})\psi_m(\mathbf{r}')\psi_n(\mathbf{r}')}{\omega - (E_m - E_n)}$$



Products of states appear $\psi_n(\mathbf{r})\psi_m(\mathbf{r})$.

Product of states in LCAO ansatz

LCAO ansatz for a molecule

$$\psi_n(\mathbf{r}) = \sum_a X_a^n f^a(\mathbf{r} - \mathbf{R}_a) \quad ,$$

$$\text{Local} \quad f^a(\mathbf{r}) = f_{l_a}^a(r) Y_{l_a, m_a}(\mathbf{r}).$$

$$\Rightarrow \quad \text{Product of states} \quad \psi_n(\mathbf{r})\psi_m(\mathbf{r}) = \sum_{a,b} X_a^n X_b^m f^a(\mathbf{r})f^b(\mathbf{r})$$

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$$H\psi_n(\mathbf{r}) = E\psi_n(\mathbf{r})$$



MBPT_LCAO

$$\chi_0(\mathbf{r}, \mathbf{r}', \omega) = \sum_{n,m} \frac{\psi_n(\mathbf{r})\psi_m(\mathbf{r})\psi_m(\mathbf{r}')\psi_n(\mathbf{r}')}{\omega - (E_m - E_n)}$$

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LCAO is too miser for products of AOs $f^a(\mathbf{r})f^b(\mathbf{r})$: we cannot express the products within the same basis $\{f^i(\mathbf{r})\}$

$$f^a(\mathbf{r})f^b(\mathbf{r}) \neq \sum_i C_i^{ab} f^i(\mathbf{r})$$



Spatial degrees of freedom: product basis ansatz

... LCAO is parsimonious, but one constructs other basis set $\{F^\mu(\mathbf{r})\}$

$$f^a(\mathbf{r})f^b(\mathbf{r}) = \sum_{\mu} V_{\mu}^{ab} F^{\mu}(\mathbf{r}).$$

All known procedures use **localized** product functions $F^{\mu}(\mathbf{r})$

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This is where **LCAO can be cheaper** than **plane waves**: $e^a e^b = e^{a+b}$

$$e^{i\mathbf{G}\mathbf{r}} e^{-i\mathbf{G}'\mathbf{r}} = V_{\mathbf{G}''}^{\mathbf{G}\mathbf{G}'} e^{i\mathbf{G}''\mathbf{r}}$$

$O(N^2)$ elements in

$$V_{\mathbf{G}''}^{\mathbf{G}\mathbf{G}'} = \delta_{\mathbf{G}-\mathbf{G}',\mathbf{G}''}$$

$$f^a(\mathbf{r})f^b(\mathbf{r}) = V_{\mu}^{ab} F^{\mu}(\mathbf{r})$$

$O(N)$ elements in V_{μ}^{ab}



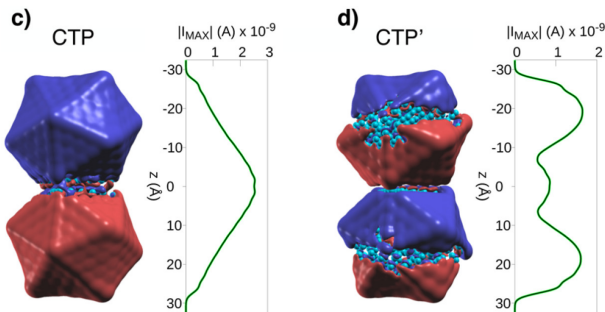
Compare with the revolutionary **FFT**: going from $O(N^2) \rightarrow O(N)$

Price to pay: implementation must use **sparse matrices**.

Product basis: (Marc Barbry yesterday)

We proposed a competitive solution for expanding density-like quantities for finite systems ¹

$$f^a(\mathbf{r})f^b(\mathbf{r}) = V_{\mu}^{ab}F^{\mu}(\mathbf{r})$$



Imaginary part of the induced density and the corresponding modulus of the electron current flowing through the middle of the dimer $2 \times \text{Na}_{380}$ ²

¹PK, FM, DF, DSP, *J. Phys: Cond. Mat.* **28** (2016) 214001.

²FM, PK, MB, JA, DSP, *ACS Photonics* **3** (2016) 269.

Closed-shell molecules belonging to G2/97 test set

<i>N</i>	Molecule name	Hill sys.	Expt.	E_{PBE}	Δ_{PBE}	GW_{PBE}	GW_{HF}
001	Borane, Trifluoro-	BF ₃	15.96	10.14	14.91	14.73	15.68
002	Borane, Trichloro-	BCl ₃	11.64	7.69	11.00	10.65	11.56
003	Aluminum Trifluoride	AlF ₃	15.45	9.81	14.00	14.14	15.03
004	Aluminum Trichloride	AlCl ₃	12.01	8.06	11.13	11.03	11.90
005	Carbon Tetrafluoride	CF ₄	14.70	10.41	14.73	14.87	16.04
006	Carbon Tetrachloride	CCl ₄	11.69	7.65	10.62	10.48	11.50
007	Carbonyl Sulfide	COS	11.19	7.57	11.62	10.52	11.16
008	Carbon Disulfide	CS ₂	10.09	6.85	10.32	9.37	9.95
009	Carbonic Difluoride	CF ₂ O	13.60	8.61	13.76	12.45	13.44
010	Silicon Tetrafluoride	F ₄ Si	16.40	10.81	14.74	15.16	16.24
011	Silane, Tetrachloro-	Cl ₄ Si	11.79	8.16	10.97	10.97	11.97
012	Nitrous Oxide	N ₂ O	12.89	8.41	13.29	11.88	12.57

ME			0.00	-4.07	0.19	-0.65	0.06
MAE			0.00	4.07	0.44	0.74	0.27

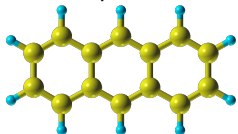
G2/97 test set contains 117 closed shell molecules

SIESTA basis set with extended AOs: Split DZP + two diffuse functions per angular momentum.

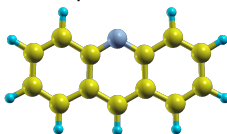
Acceptor molecules: reference

G_0W_0 done with FHI-aims package, all-electron starting point³

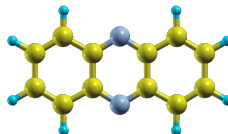
For 24 acceptor molecules. Examples



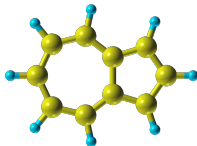
Anthracene



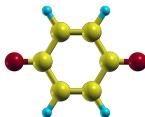
Acridine



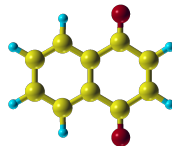
Phenazine



Azulene



Benzenoquinone



Naphtalenoquinone

³JWK, XW, LG *et al* J. Chem. Theo. Comput. 12 (2016) 615-626

Acceptor molecules: basis sets for them (Federico's talk!)

- ▶ Very large basis sets were employed in FHI-aims,

Tier4: [min]+4s4p4d3f2g

- ▶ We use correlation-consistent cc-pVXZ basis sets

<https://bse.pnl.gov/bse/portal>

core shells replaced by pseudo-potentials;

occupied valence shells are replaced by solutions of
radial Kohn-Sham equation ([min] in SIESTA)

Example of taken Gaussian exponents α_p
for cc-pVTZ (X=3), carbon (Z=6)

$$f^o(\mathbf{r}) = e^{-\alpha_p r^2} Y_{j_p m}(\mathbf{r})$$

Z	X	j_p	pAO	$-\alpha_p$
6	3	0	1	0.9059
6	3	0	2	0.1285
6	3	1	3	0.3827
6	3	1	4	0.1209
6	3	2	5	1.097
6	3	2	6	0.318
6	3	3	7	0.761

Acceptor molecules: Hartree-Fock, –HOMO

Molecule	AIMS	2ZP	3ZP	3ZP,f	4ZP	4ZP,f	5ZP	5ZP,f
anthracene	7.0643	7.4159	7.1648	7.1861	7.1064	7.1175	7.0792	7.0854
acridine	7.6708	8.0280	7.7730	7.7957	7.7168	7.7281	7.6887	7.6950
phenazine	8.1565	8.5275	8.2643	8.2869	8.2068	8.2181	8.1772	8.1836
azulene	7.1062	7.4293	7.1906	7.2149	7.1379	7.1497	7.1134	7.1196
BQ	11.224	11.559	11.305	11.329	11.244	11.257	11.215	11.223
NQ	9.8206	10.171	9.9095	9.9368	9.8530	9.8662	9.8241	9.8314
				
ME	0.0000	0.4031	0.0963	0.1147	0.0262	0.0350	0.0068	0.0020
MAE	0.0000	0.4031	0.0963	0.1147	0.0320	0.0400	0.0138	0.0135
Max	0.0000	0.6094	0.1766	0.1753	0.0530	0.0616	0.0740	0.0677

HF: a small discrepancy $\text{MAE} = 0.0135 \text{ eV}$

Acceptor molecules: $\text{HF}+G_0W_0$, $-\text{HOMO}$

Molecule	AIMS	2ZP	3ZP	3ZP,f	4ZP	4ZP,f	5ZP	5ZP,f
anthracene	7.721	7.587	7.514	7.650	7.513	7.618	7.508	7.581
acridine	8.246	8.109	8.028	8.171	8.026	8.135	8.020	8.096
phenazine	8.682	8.559	8.479	8.624	8.474	8.585	8.467	8.543
azulene	7.672	7.538	7.488	7.628	7.493	7.598	7.490	7.562
BQ	11.27	11.58	11.56	11.58	11.49	11.54	11.37	11.51
NQ	10.23	10.13	9.993	10.14	10.01	10.10	9.987	10.06
...								
ME	0.000	0.007	-0.150	-0.010	-0.172	-0.091	-0.186	-0.130
MAE	0.000	0.124	0.198	0.076	0.213	0.147	0.218	0.181
Max	0.000	0.474	0.296	0.434	0.288	0.416	0.324	0.494

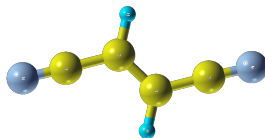
$\text{HF}+G_0W_0$: a small but sensible discrepancy $\text{MAE} = 0.181$ eV.

Similar error in electron affinities $\text{MAE} = 0.241$ eV.

Absence of core electrons cause such discrepancy (≈ 0.18 eV) ?

Influence of core electrons

- ▶ All electron calculation with cc-pVDZ for Fumaronitrile



- ▶ Core electrons can influence Σ_c via G $\Sigma_c = \sum_n \frac{|\psi_n\rangle\langle\psi_n|}{\omega - i\omega' - E_n} W_c(i\omega')$
Core electrons from G is negligible

- ▶ Core electrons can influence Σ_c via χ $\Sigma_c = G(1 - v\chi)^{-1}v\chi v$

Contour def.	low + high	$\omega_{nm} = E_n - E_m$	
		low	high
\int	-1.0504	-1.0035	-0.0470
Res	0.9818	0.9803	0.00153
$\int + \text{Res}$	-0.0686	-0.0232	-0.04546



- ▶ Core electrons influence significantly the GW correction!

Conclusions & Problem to be solved

- ▶ $G_0 W_0$ can be computationally cheaper than Hartree-Fock (diagonal Σ_c , contour deformation, HOMO-LUMO only)
- ▶ Absence of core electrons causes a small, but *sensible* error.
- ▶ Small basis sets (3ZP,f) can be rather precise!
- ▶ Core-polarization potential (CPP) must be realized in our framework

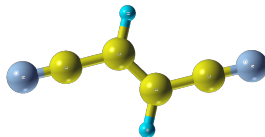
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Thank you for your kind attention!

Calculations for Fumaronitrile

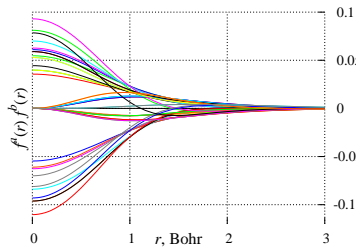
- All electron calculation with cc-pVDZ for Fumaronitrile



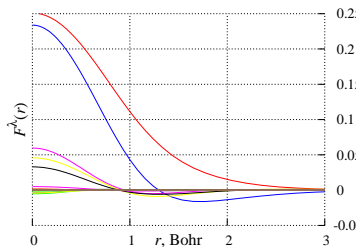
Molecule	AIMS	2ZP	3ZP	3ZP,f	4ZP	4ZP,f	5ZP	5ZP,f
Fumaronitrile	11.833	11.801	11.650	11.787	11.600	11.710	11.580	11.658

Implications of products of orbitals

- There are a lot of collinearity in products of orbitals $f^a(\mathbf{r})f^b(\mathbf{r})$



- However, there are **linear combinations** of products that would be mutually orthogonal

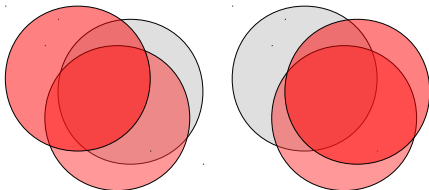


Dominant products

- Use the **Coulomb metric**

$$g^{ab,cd} \equiv \iint \frac{f^a(\mathbf{r})f^b(\mathbf{r})f^c(\mathbf{r}')f^d(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|} d^3r d^3r'.$$

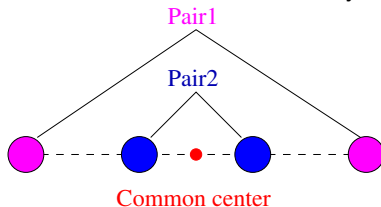
- A global metric $g^{ab,cd}$ of a big molecule is too big \Rightarrow **lets diagonalize for each pair of atoms individually**⁴



⁴D. Foerster, J. Chem. Phys. **128**, 034108 (2008).

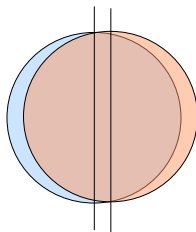
Dominant products: contra

- ▶ DP is too big comparing to PW...
- ▶ DP generates too overlapping functions
- ▶ DP can put extra functions where there are already functions

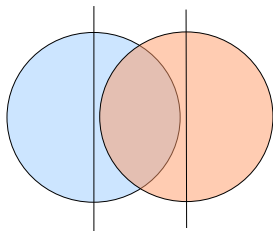


Atom-centered vertex: geometry problem

- ▶ Nuclei separations are normally smaller than orbital's spatial support

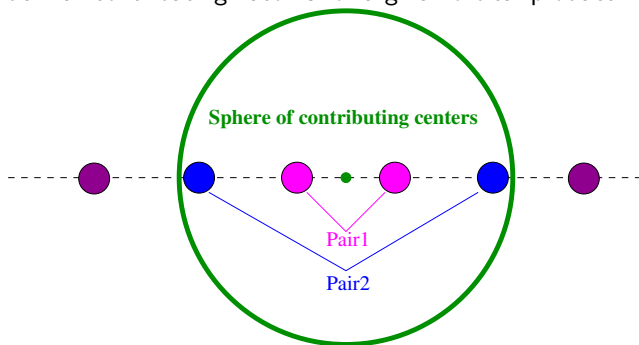


- ▶ If nuclei separations are comparable to the support then this product is small and not important. . .



Atom-centered vertex: geometry solution

- ▶ Method: put some basis functions exclusively at atom nuclei and define “contributing” atoms for a given orbital product.



- ▶ Corollary: functions get reused for **Pair1** and **Pair2**.

Atom-centered vertex: contra and how to deal with it

- ▶ For atom-centered product basis we have $f^a(\mathbf{r})f^b(\mathbf{r}) = V_{\mu}^{ab}A^{\mu}(\mathbf{r})$
- ▶ For dominant products we have $f^a(\mathbf{r})f^b(\mathbf{r}) = V_{\tilde{\mu}}^{ab}F^{\tilde{\mu}}(\mathbf{r})$
- ▶ $N(A^{\mu}) \approx 1/4 \cdot N(F^{\tilde{\mu}})$ 😊
- ▶ but $N(V_{\mu}^{ab}) \approx 10 \cdot N(V_{\tilde{\mu}}^{ab})$ 😞 (Fock exchange 100x slower!)
- ▶ However, it is enormously more advantageous to factorize

$$V_{\mu}^{ab} = V_{\tilde{\mu}}^{ab} c_{\mu}^{\tilde{\mu}}$$

i.e.

$$\text{Atom-centered} = \text{Dominant} \times \text{Conversion}$$

Atom-centered vertex: algebra

- ▶ We have dominant products $f^a(\mathbf{r})f^b(\mathbf{r}) = V_{\tilde{\mu}}^{ab}F^{\tilde{\mu}}(\mathbf{r})$
- ▶ Let's express $F^{\tilde{\mu}}(\mathbf{r})$ in terms of atom-centered functions $A^{\mu}(\mathbf{r})$

$$F^{\tilde{\mu}}(\mathbf{r}) = c_{\mu}^{\tilde{\mu}}A^{\mu}(\mathbf{r})$$

Coulomb metric works best \Rightarrow

$$\int F^{\tilde{\mu}}(\mathbf{r}) \frac{1}{|\mathbf{r} - \mathbf{r}'|} A^{\nu}(\mathbf{r}') d\mathbf{r} d\mathbf{r}' = c_{\mu}^{\tilde{\mu}} \int A^{\mu}(\mathbf{r}) \frac{1}{|\mathbf{r} - \mathbf{r}'|} A^{\nu}(\mathbf{r}') d\mathbf{r} d\mathbf{r}'$$

- ▶ Invert square matrix $V^{\mu\nu}$:

$$c_{\mu}^{\tilde{\mu}} = R^{\tilde{\mu}\nu} V_{\nu\mu}^{-1}.$$

$O(N^4)$ response function

$$\chi_{\mu\nu}^0(\omega) = 0$$

Loop over $\{mn\}$ chunks

$$A_{\tilde{\mu}}^{an} = V_{\tilde{\mu}}^{ab} X_b^n$$

$$B_{\tilde{\mu}}^{mn} = X_a^m A_{\tilde{\mu}}^{an}$$

$$B_{\mu}^{mn} = B_{\tilde{\mu}}^{mn} c_{\mu}^{\tilde{\mu}}$$

Loop over ω

$$C_{\mu}^{mn} = B_{\mu}^{mn} \frac{f_n - f_m}{(\omega - (E_m - E_n))}$$

$$\chi_{\mu\nu}^0(\omega) = \chi_{\mu\nu}^0(\omega) + \sum_{mn} B_{\mu}^{mn} C_{\nu}^{mn}.$$

The algorithm to compute the response function for atom-centered product basis using sum-over-states expression. Critical part of the algorithm is the computation of the molecular-orbital product vertex B_{μ}^{mn} . It is computed utilizing the factorization of the atom-centered product vertex according to equation

$$V = \tilde{V} \cdot c$$

This is much faster than time-frequency domain techniques...