

NanoSpain Conference 2010, Malaga (Spain)

Electronic excitations in thin metallic films

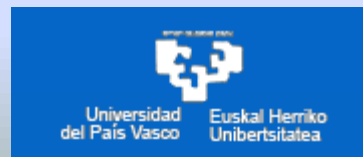
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1. Introduction and motivation

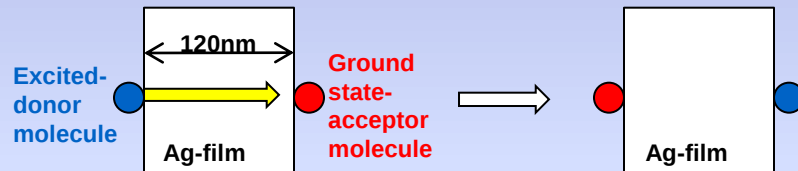
Why is proper knowledge of the electronic excitations at metallic surface important?

1. Plasmonic circuits-guided surface plasmons (polariton) could be subject to significant Landau damping.

Berni P. et al, J. Appl. Phys. **98** 043109 (2005)

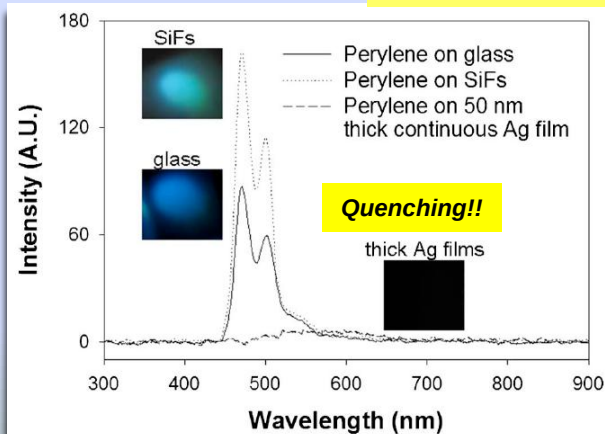
2. Plasmonic devices-thin silver film covered from both sides by thin polymer films with molecular chromophores. **ENERGY TRANSFER ACROSS A METAL FILM MEDIATED BY SURFACE PLASMON POLARITONS**

Andrew P. and Barnes W. L., Science **306** 1002 (2004)



3. Biosensing –Metal-enhanced fluorescence: Surface plasmon can radiate a Fluorophore's structured emission.

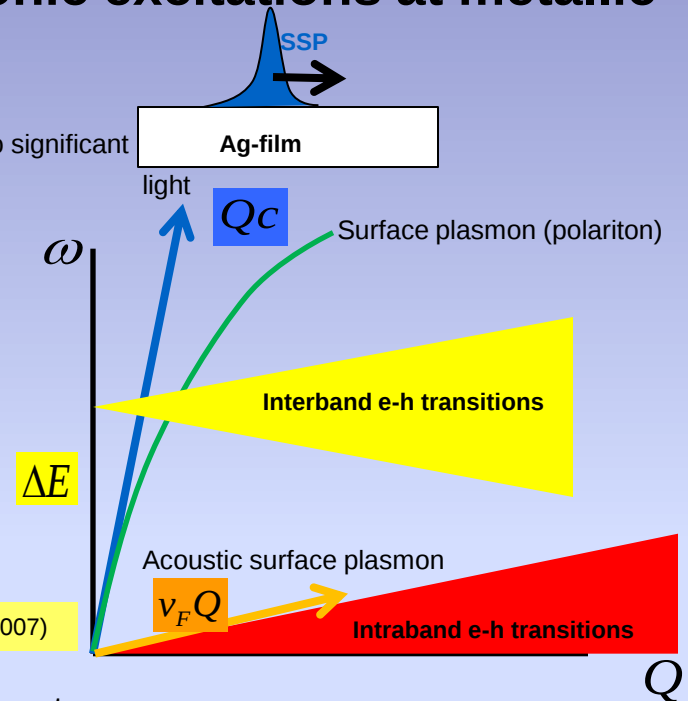
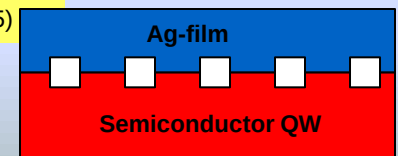
Yongxia Zhang, et al, Appl. Phys. Lett. **90**,053107 (2007)



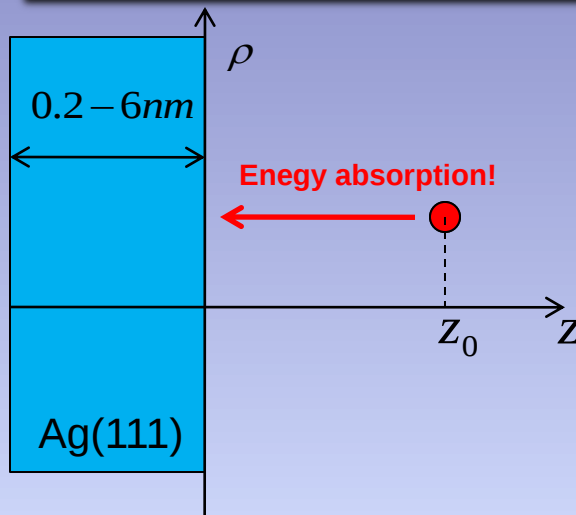
4. Light-emitting structures semiconductor QW coated by thin silver films increase spontaneous emission in QW. Efficient energy transfer from electron-hole pair recombination (in QW) to a surface plasmon (in silver film). Energy is converted into light by using corrugated metal-semiconductor interface.

A. Neogi, et al, Phys. Rev. **B 66**, 153305 (2002)

K. Okamoto, et al, Appl. Phys. Lett. **87** 071102 (2005)



2. Mathematical formulation and modeling of the system



Sudden created point charge near Ag(111) surface,
V.M. Silkin, Poster **90**

$$V^{ext} = \frac{2\pi e}{Q} e^{-Q|z-z_0|} \cos(\omega t)$$

$$P_{loss}(Q, \omega) \propto \text{Im} W^{ind}(Q, \omega, z_0, z_0)$$

$$W^{ind}(Q, \omega_0, z_0, z_0) = e^{-2Qz_0} g(Q, \omega)$$

Surface response function

$$g(Q, \omega) = \int dz_1 dz_2 e^{2Q(z_1+z_2)} R(Q, \omega, z_1, z_2)$$

$-\text{Im} g(Q, \omega)$ Intensity of the real excitations in the system!

Response function-RPA

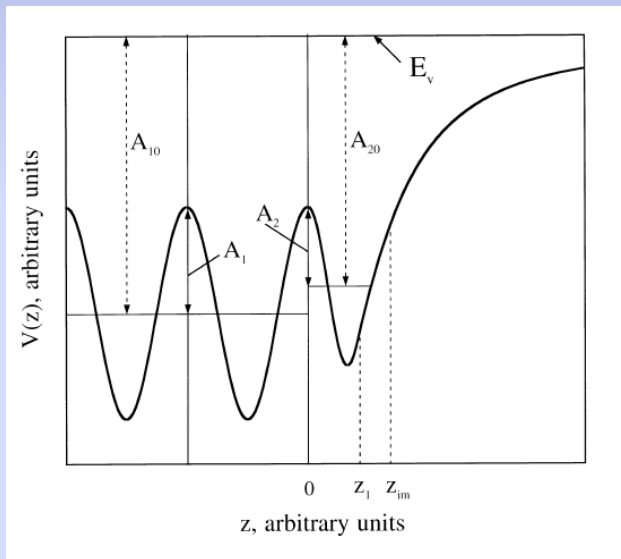
$$R = \frac{R_0}{1 - v R_0}$$

Free electron response function

$$R^0(\mathbf{Q}, \omega, z, z') = \sum_{n,m,\mathbf{K}} \frac{f(E_{n,\mathbf{K}}) - f(E_{m,\mathbf{K}+\mathbf{Q}})}{\omega - (E_{m,\mathbf{K}+\mathbf{Q}} - E_{n,\mathbf{K}}) + i\eta} \Phi_n(z) \Phi_m(z) \Phi_n(z') \Phi_m(z')$$

Band structure and wave function-1D model potential

E.V. Chulkov, V. M. Silkin and P.M. Echenique, Surf.Sci. **437** (1999) 330-352



Parameters of the potential are set up to fit experimental (first principle) datas for surface state, first image state and postion and width of band gap at $\bar{\Gamma}$ point

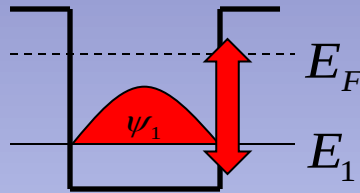
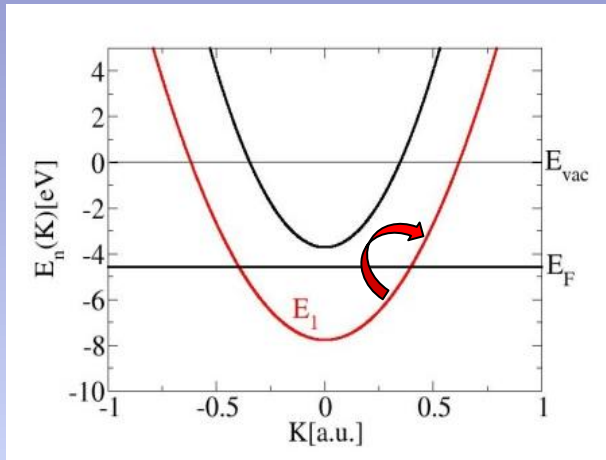
Band structure and wave functions- solutions of 1D Schrodinger equation:

$$\left[-\frac{\hbar^2}{2m^*} \frac{d^2}{dz^2} + V_{MP}(z) \right] \Phi_n(z) = E_n \Phi_n(z)$$

$$E_{n,\vec{K}} = E_n + \frac{\hbar^2 K^2}{2m^*}$$

3. Excitation spectra inside ultra thin silver films

1ML(0.2nm)- Ag(111)



1ML-one quasi 2D electron gas

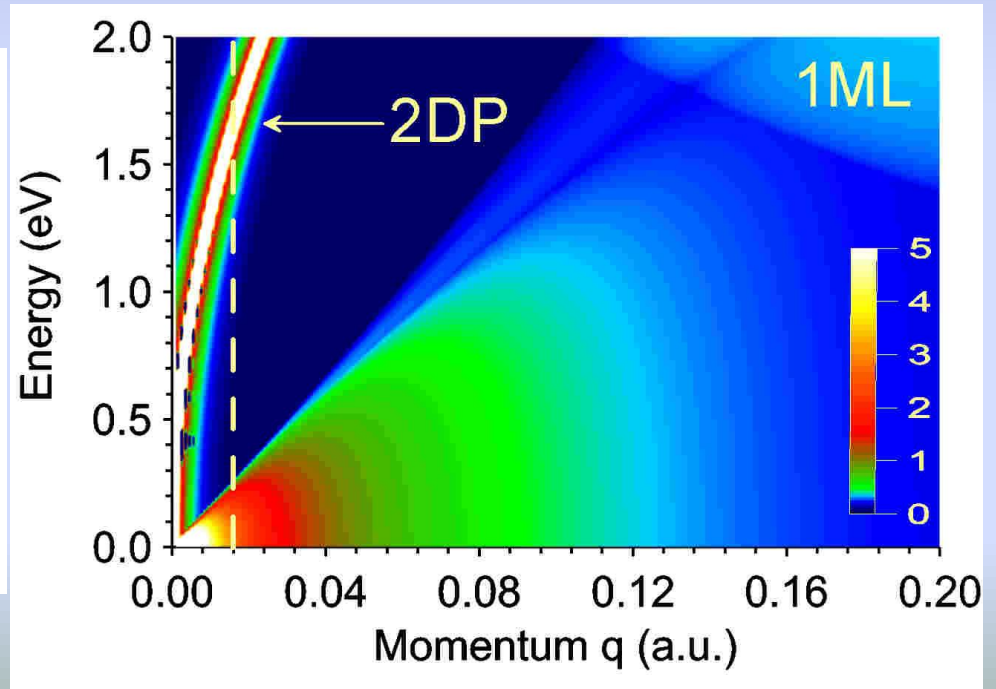
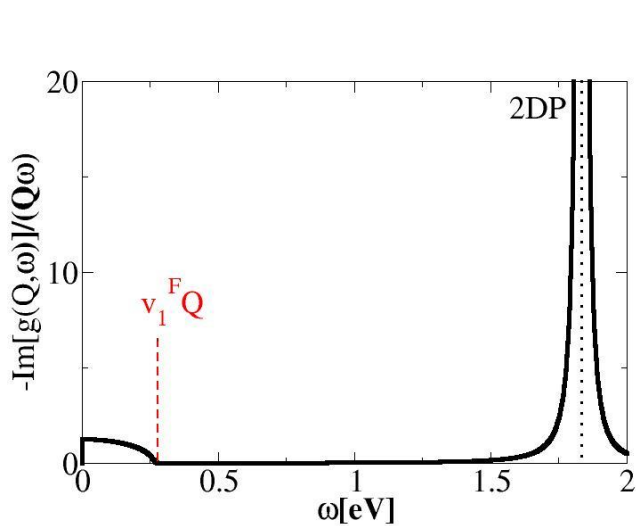
$$\epsilon_{eff}^{2D}(Q, \omega) = 1 - V_Q R_1^{2D}(Q, \omega)$$

$$V_Q = v_Q \int_{-\infty}^{\infty} dz_1 dz_2 \psi_1^2(z_1) e^{-Q(z_1 - z_2)} \psi_1^2(z_2) \xrightarrow{Q \rightarrow 0} \frac{2\pi e}{Q}$$

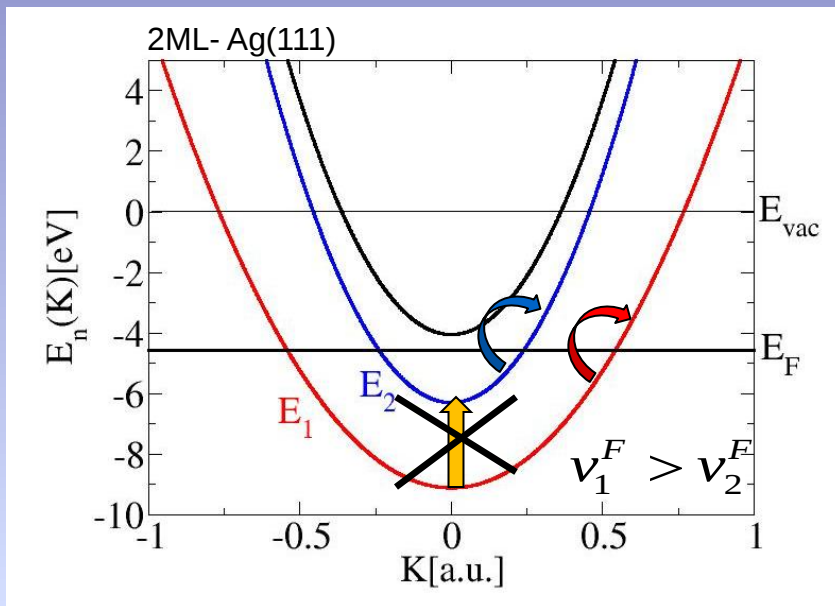
$$\text{Im} \epsilon_{eff}(Q, \omega) = 0$$

$$\text{Re} \epsilon_{eff}(Q, \omega) = 0$$

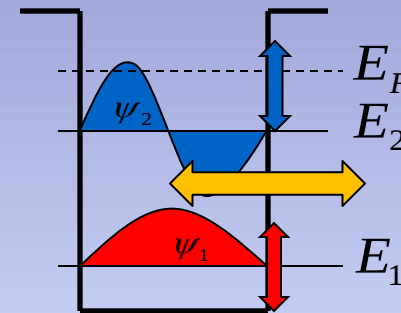
$$\omega_{2D} = \frac{k_1^F}{\sqrt{m^*}} \sqrt{Q}$$



2ML(0.4nm)- Ag(111) film



2ML-two quasi 2D electron gases



Approximations:

a) Longwavelength limit $Q \rightarrow 0$

b) There are no transitions between bands E_1 and E_2

$$\varepsilon_{eff}^{2D}(Q, \omega) = 1 - v_Q \left[R_1^{2D}(Q, \omega) + R_2^{2D}(Q, \omega) \right]$$

$$\omega_{ASP} = v_S Q$$

$$v_S = \frac{2}{\sqrt{3}} v_1^F < v_2^F$$

$$v_{1,2}^F = \sqrt{\frac{2}{m_{1,2}^*} (E_F - E_{1,2})}$$

2ML-Two component quasi 2D electron

plasma-Existence of fast and slow electron plasmas

allow forming of Acoustic Surface Plasmon (ASP) . D. Pines, Can J. Phys. **34**, 1379 (1956)

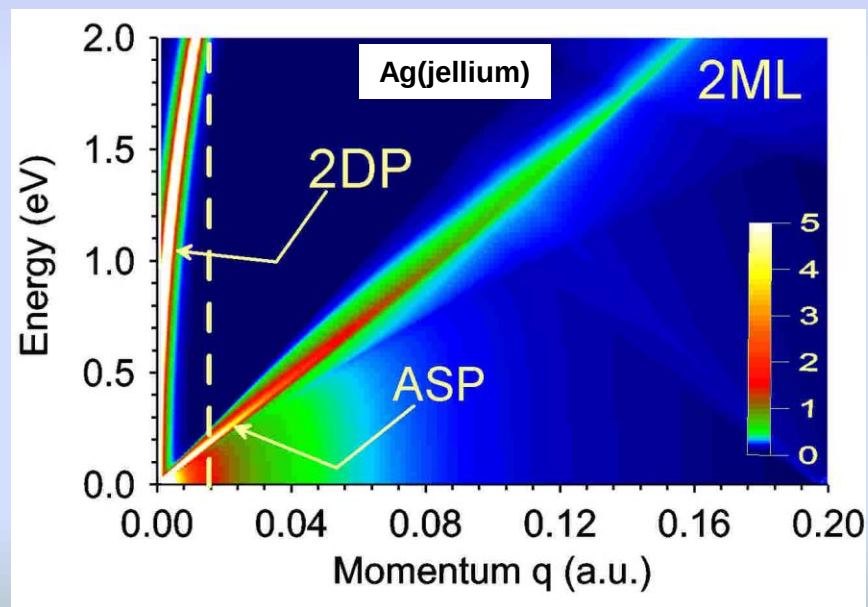
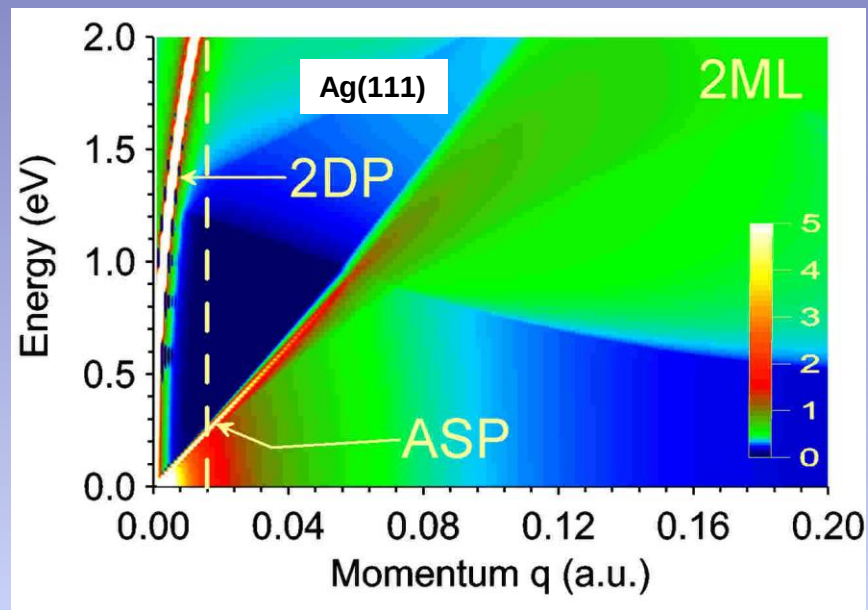
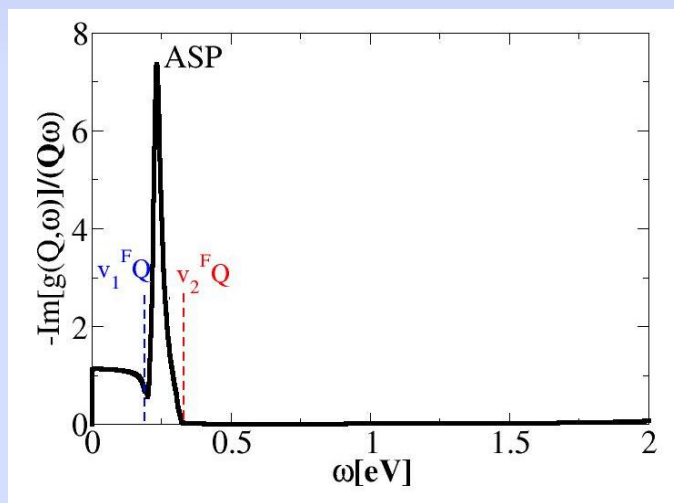
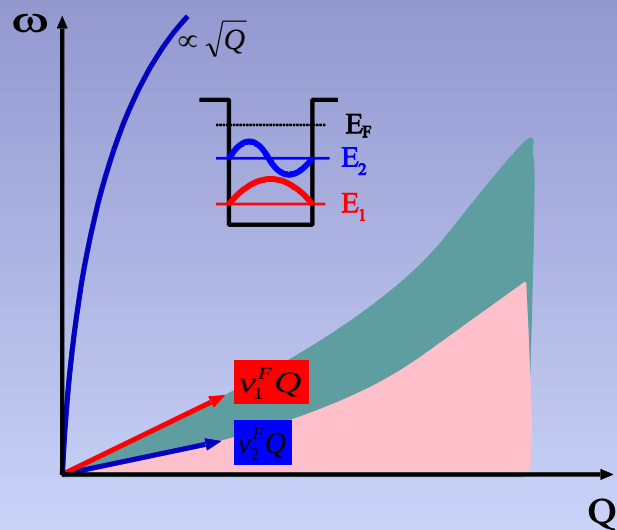
3ML, 4ML, . . . Multiplication of acoustic

surface plasmon branches

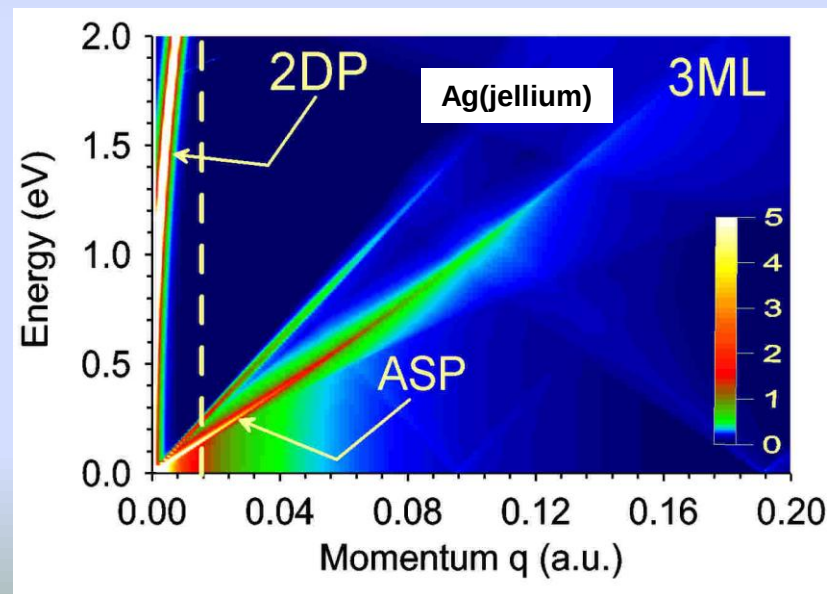
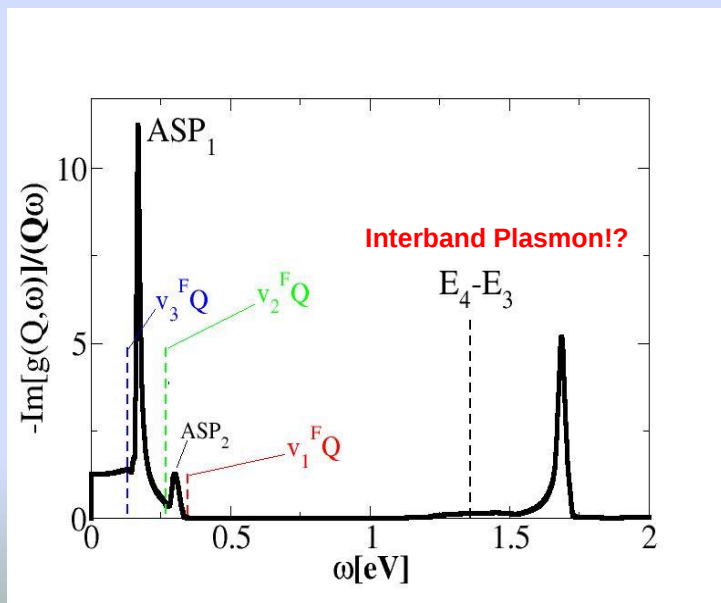
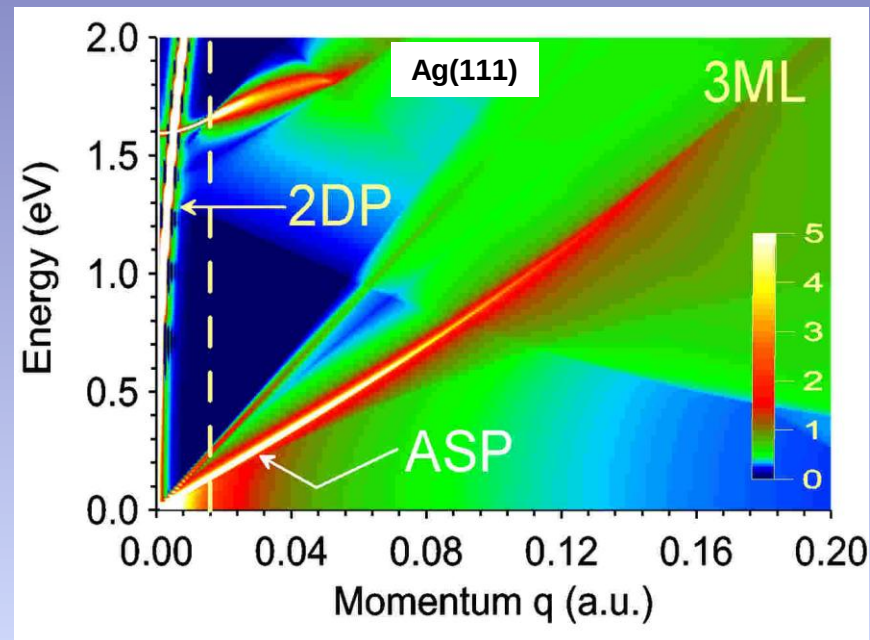
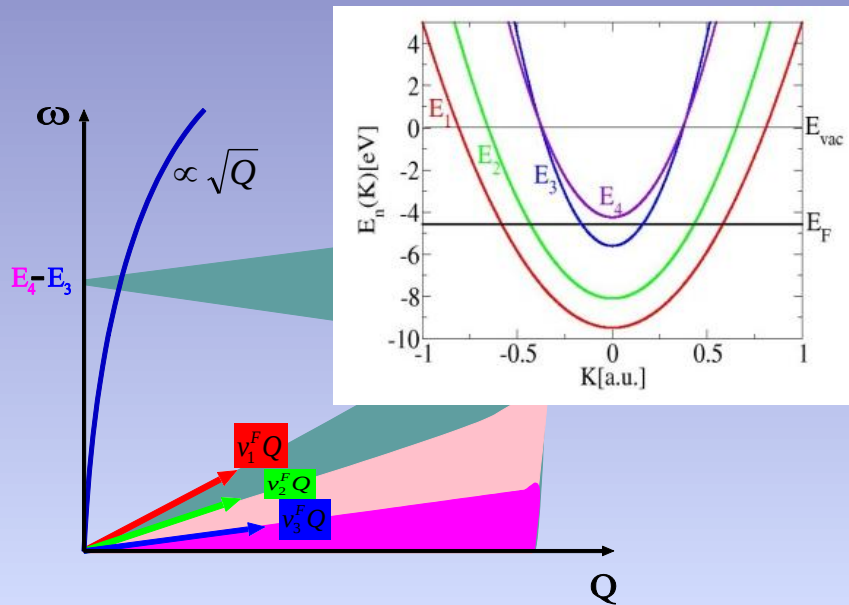
$$\omega_{ASP}^i = v_S^i Q$$

$$v_1^F < v_S^1 < v_2^F < v_S^2 < v_2^F < \dots$$

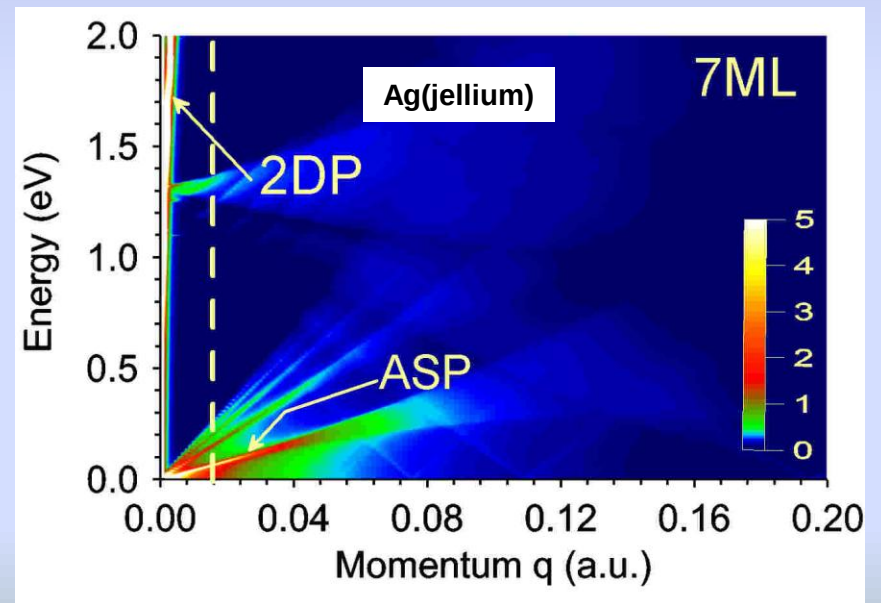
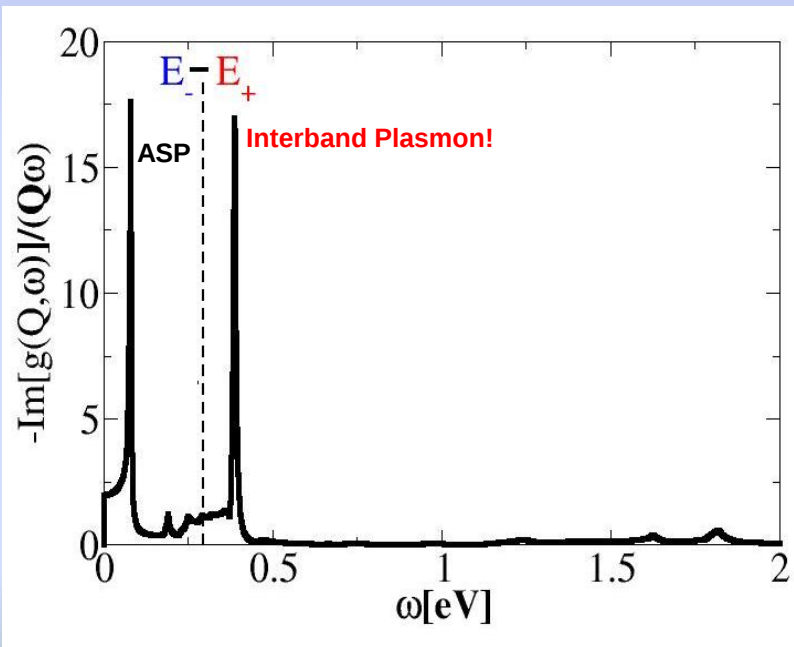
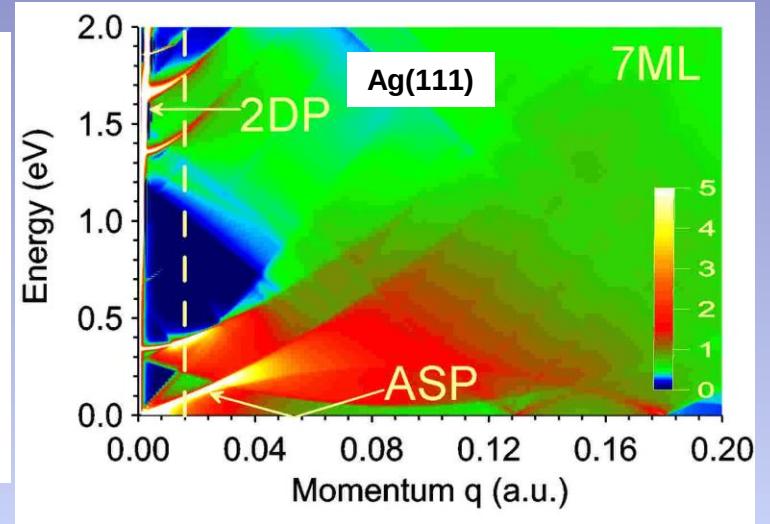
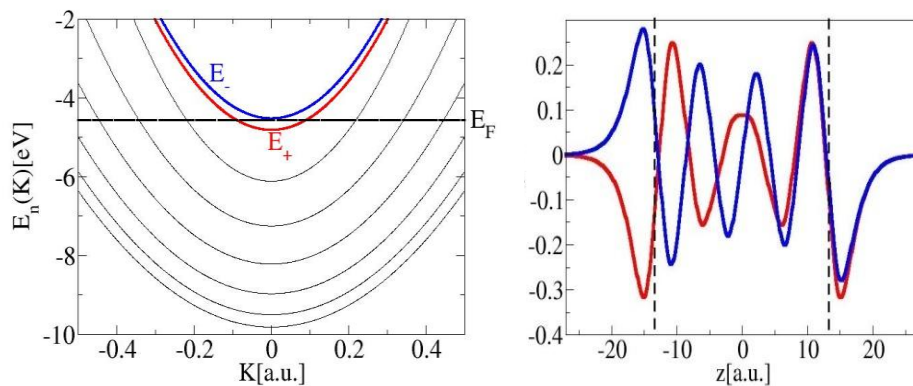
2ML(0.4nm)- Ag(111) film



3ML(0.6nm)- Ag(111) film

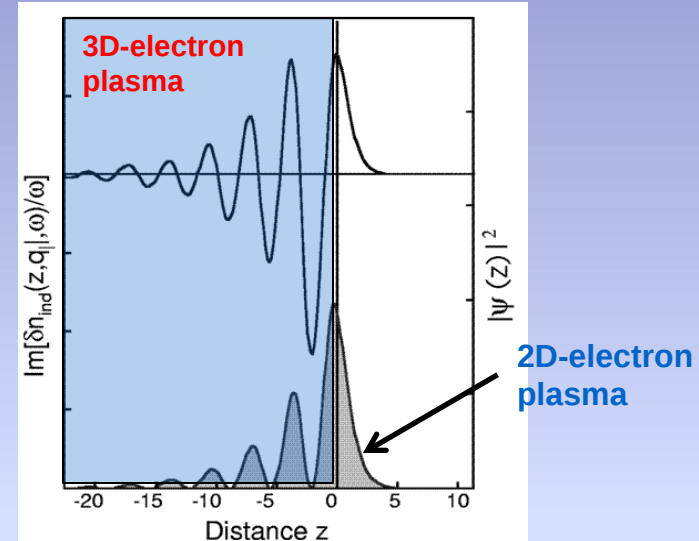
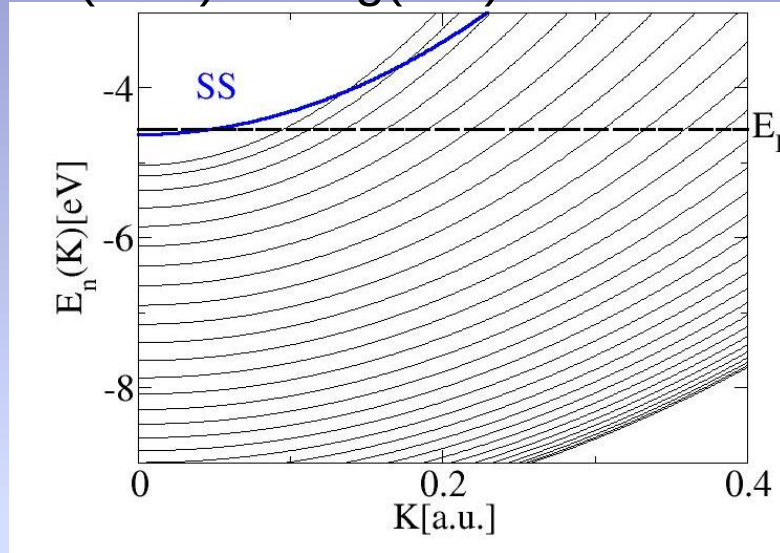


7ML(1.4nm)- Ag(111) film



4. Excitation spectra inside thick silver films and transition to spectra in semiinfinite metals

31(6nm)ML- Ag(111) film



J. M. Pitarke, V. U. Nazarov, V. M. Silkin, E. V. Chulkov, E. Zaremba, and P. M. Echenique Phys. Rev. **B 70**, 205403 (2004)

V. M. Silkin, J.M. Pitarke, E. V. Chulkov and P. M. Echenique Phys. Rev. **B 72**, 115435 (2005)

$$\varepsilon_{\text{eff}}^{2D}(Q, \omega) = 1 - W^{3D}(Q, \omega) R_{SS}^{2D}(Q, \omega)$$

Dispersion relation equation

$$\varepsilon_{\text{eff}}^{2D}(Q, \omega) = 0$$

$$\omega_{ASP} = \alpha v_{SS}^F Q$$

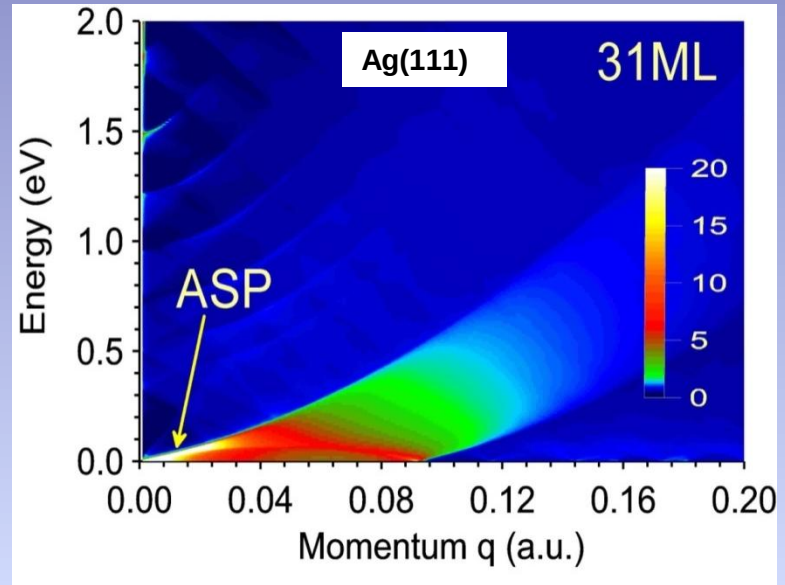
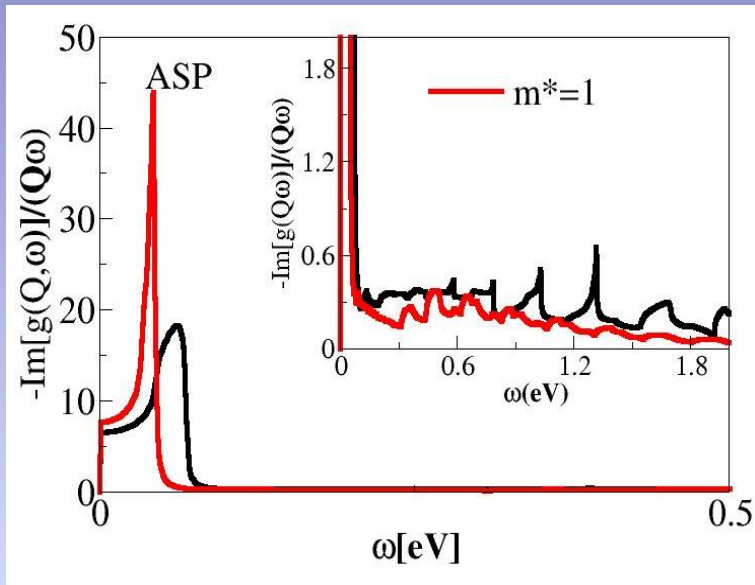
$$\alpha \approx 1$$

$$W^{3D}(Q, \omega) = \int_{-\infty}^{\infty} dz_1 dz_2 \psi_{SS}^2(z_1) W^{3D}(Q, \omega, z_1, z_2) \psi_{SS}^2(z_2)$$

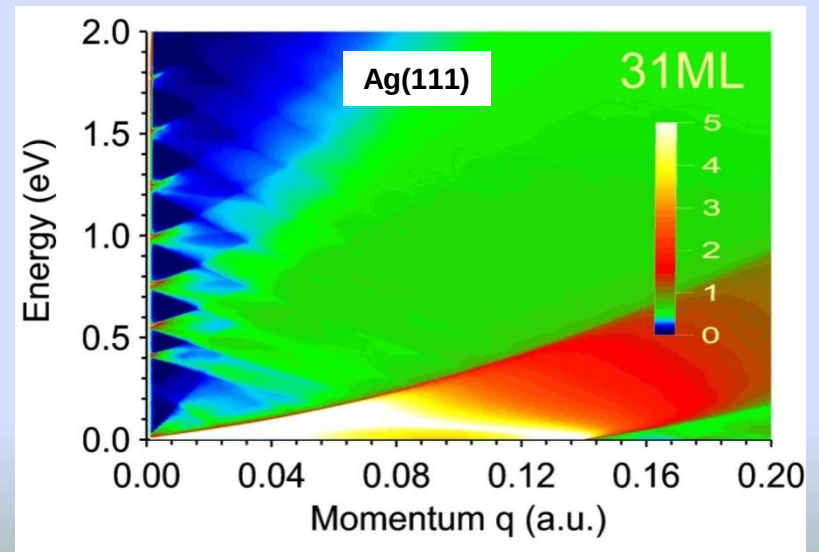
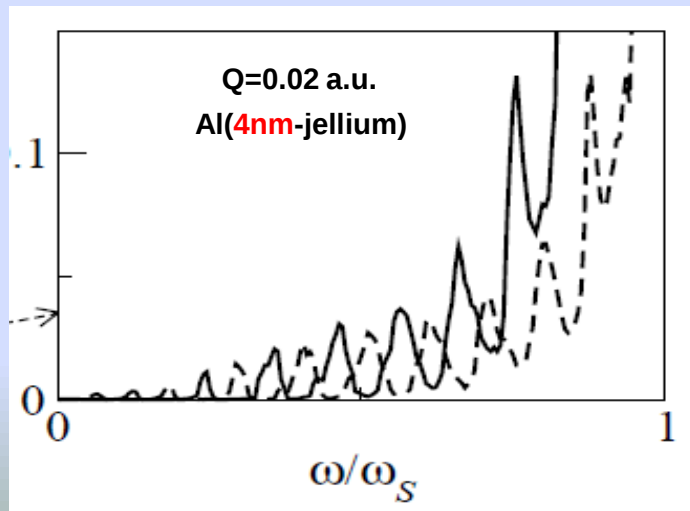
Without 3D electron plasma $W^{3D}(Q, \omega) \rightarrow v_Q = \frac{2\pi e}{Q} \Leftrightarrow$ One quasi 2D electron gas

$$\omega_{2D} \propto \sqrt{Q}$$

31(6nm)ML- Ag(111) film



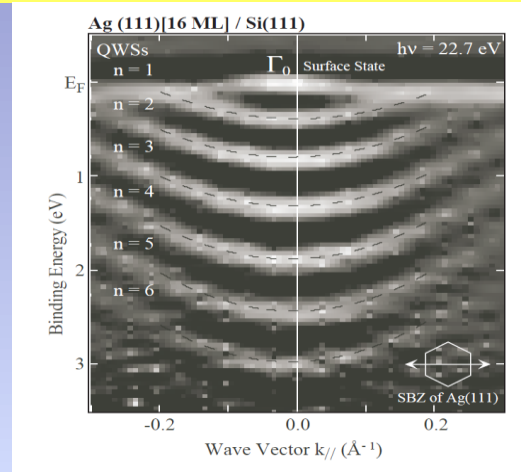
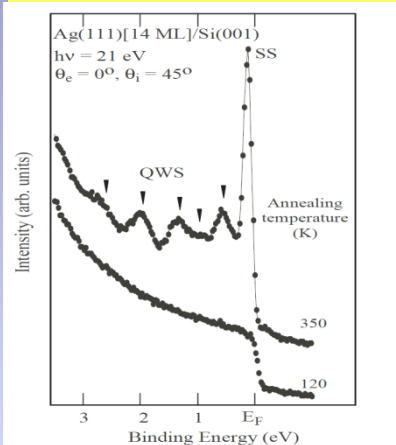
L. Marušić, V. Despoja and M. Šunjić, J. Phys.: Condens. Matter **18** (2006) 4252



5. Conclusion

ARPES measurements of quantum size effects

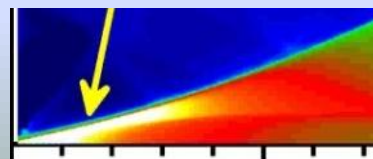
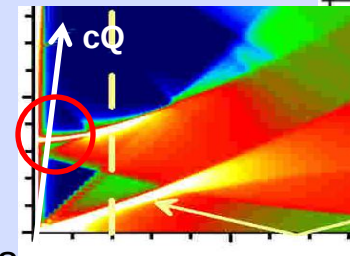
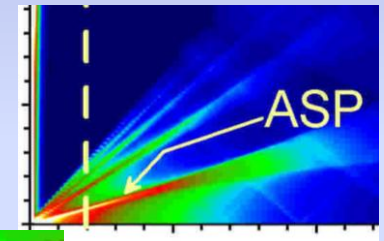
I. Matsuda *et al*, J. Surf. Sci. Nanotech. Vol. 2 (2004) 169-177



1. Ultra thin metallic films $\rightarrow$ quantum size effect $\rightarrow$ multicomponent 2D plasma $\rightarrow$ multiplication of ASP branches.

2. Thin metallic slabs $\rightarrow$ interband plasmon.
 It would be nice to investigate how it couples to light.

3. Semiinfinite noble metals (111) surfaces.
 Slow 2D plasma (surface state band) + fast 3D plasma (bulk state bands) = Acoustic Surface Plasmon



**Thank you for your
attention!**