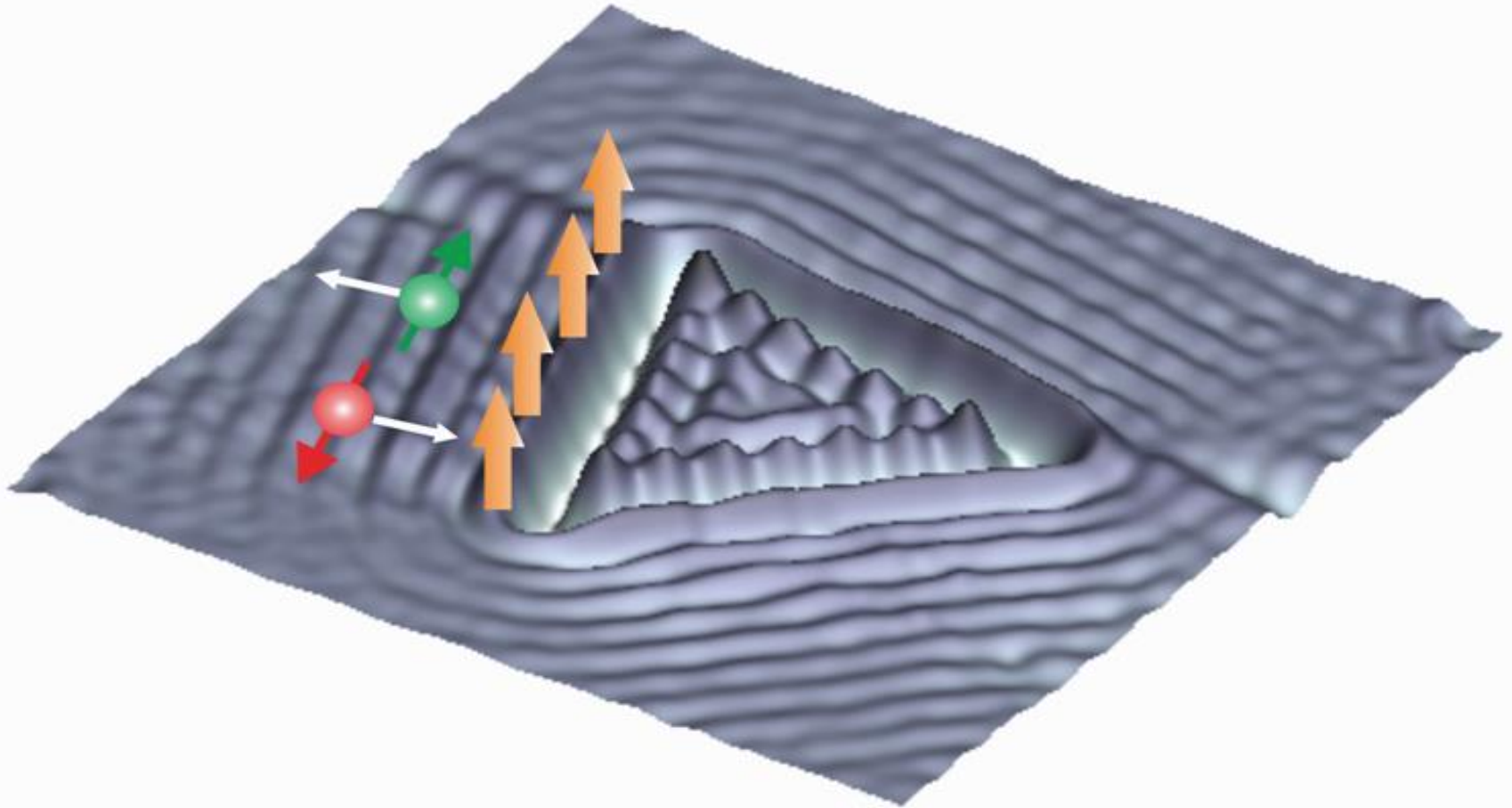


Spin-dependent electron scattering in 2D electron systems

Aitor Mugarza

Catalan Institute of Nanoscience and Nanotechnology (ICN2)

Barcelona, Spain



Generalitat de Catalunya
**Departament d'Economia
i Coneixement**



GOBIERNO
DE ESPAÑA

MINISTERIO
DE ECONOMÍA
Y COMPETITIVIDAD



ICREA



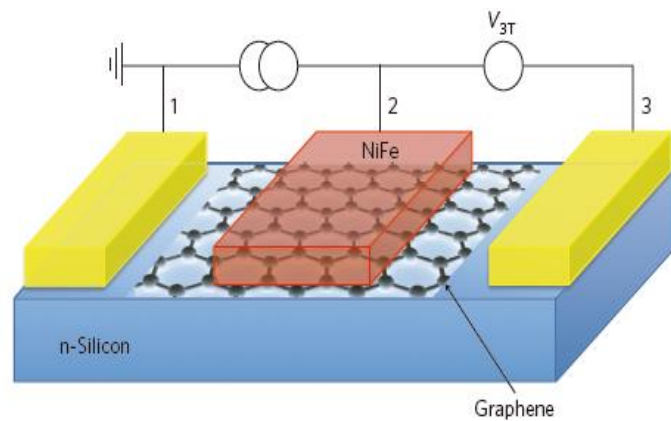
Institut Català
de Nanociència
i Nanotecnologia



EXCELENCIA
SEVERO
OCHOA

Nanospain, Donostia, March 9th 2017

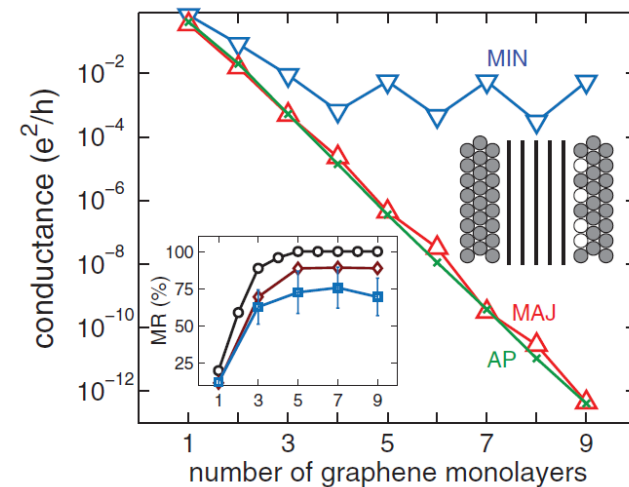
Spin filtering at the G/Ni(111) interface



In FM/semiconductor junctions:

- ☐ High injection efficiency
- ☐ Low barrier resistance

van t'Erve et al. Nat. Nanotech. 7 (2012)

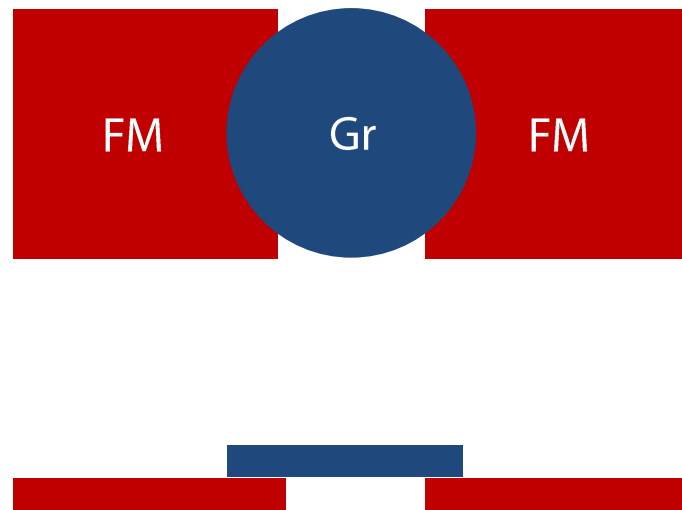


Spin filtering in G/Ni interfaces:

- ☐ Magnetoresistance close to 100%

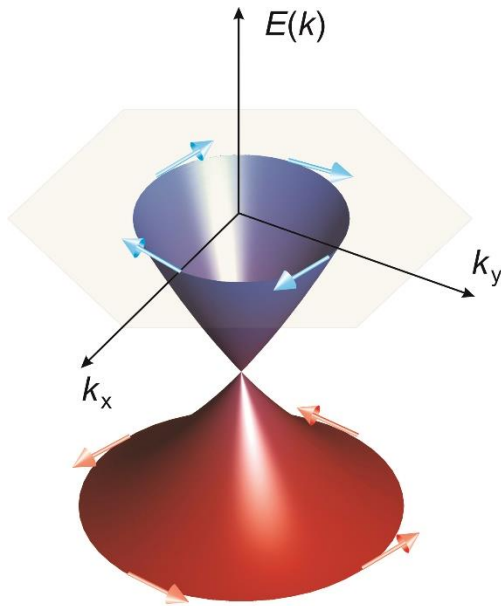
V. M. Karpan et al. PRL 99 (2007)

Spin in scattering: Exchange split states



Spin in scattering: spin-orbit coupling

Topological states



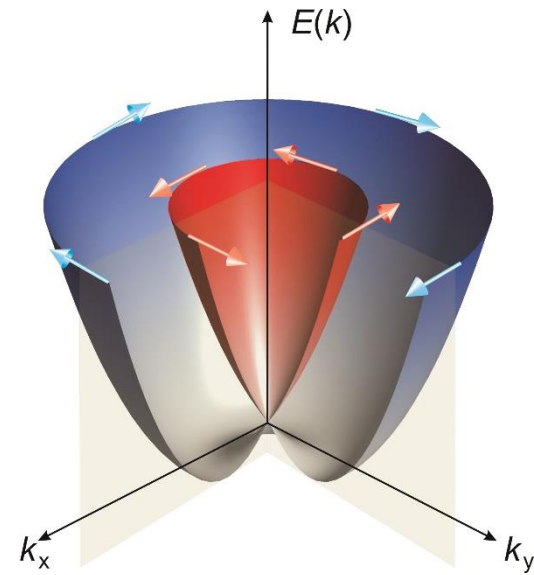
Spin locked to momentum



$$\mathbf{J}(\mathbf{x}) \leftrightarrow \mathbf{S}(\mathbf{x})$$

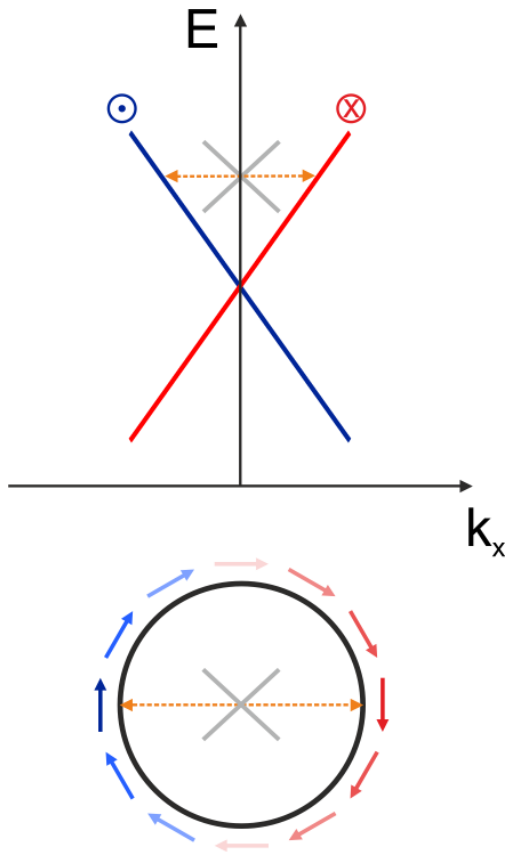
Charge current Spin density

Rashba states



Spin in scattering: spin-orbit coupling

Topological states



Spin locked to momentum



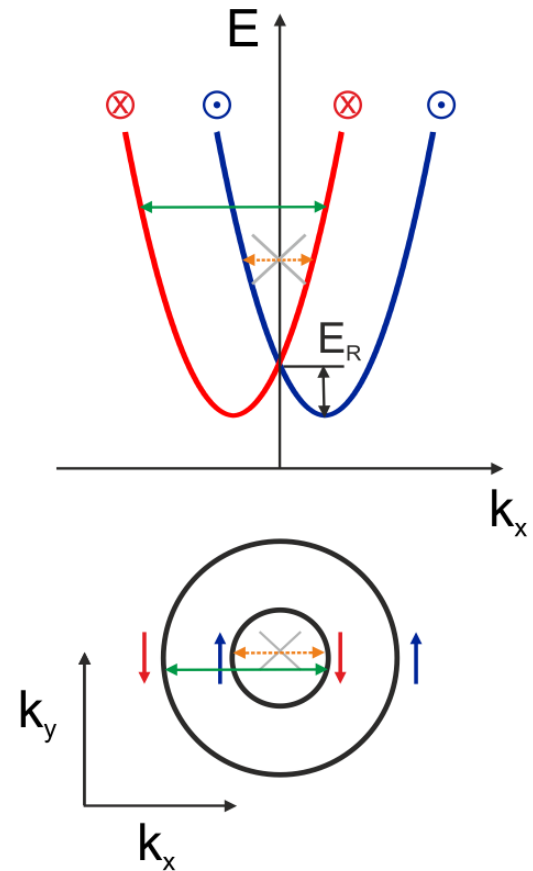
$$\mathbf{J}(\mathbf{x}) \leftrightarrow \mathbf{S}(\mathbf{x})$$

Charge
current

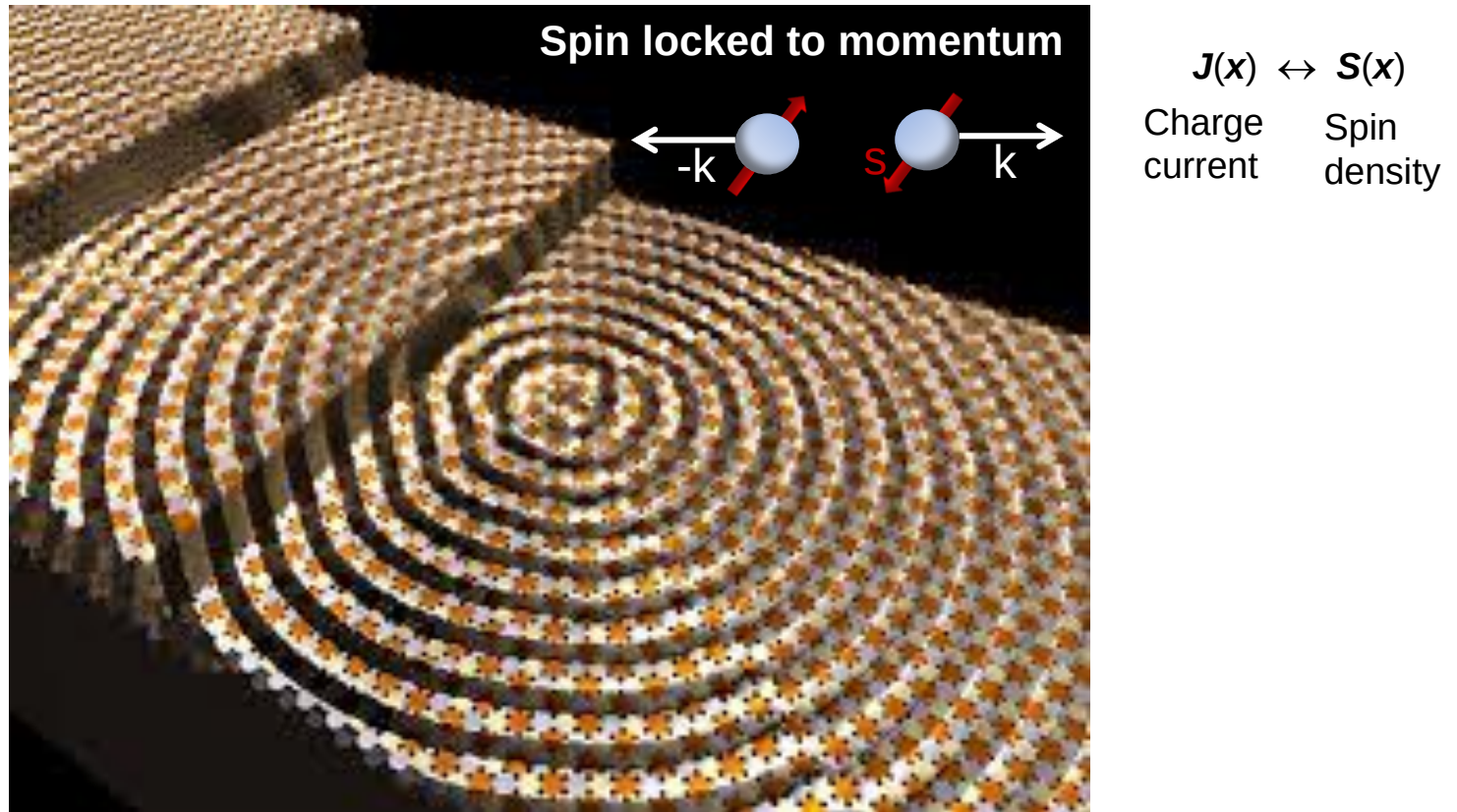
Spin
density

Quenching of backscattering
Non-dissipative currents?

Rashba states



Quenching of backscattering in topological insulators

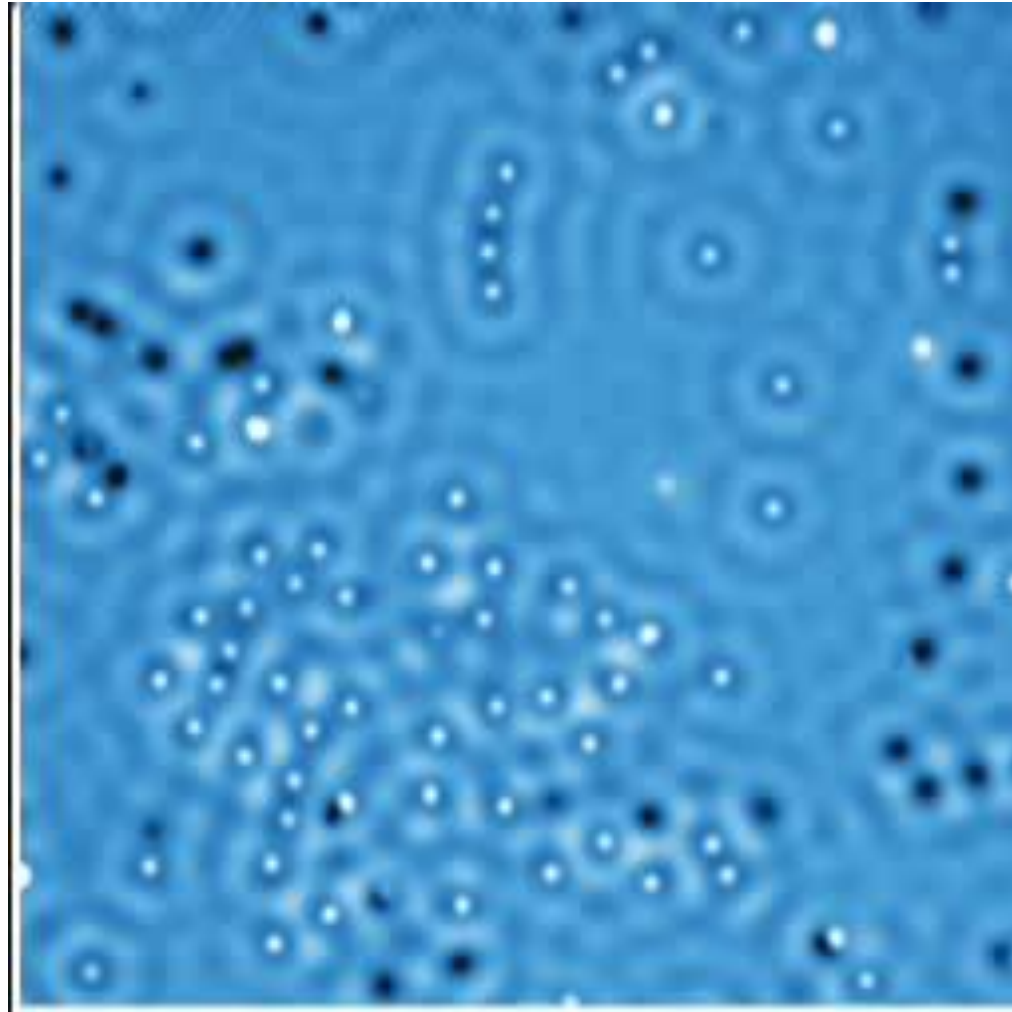


J. Seo et al. Nature **466**, 343 (2010)

Non-dissipative spin polarized currents?

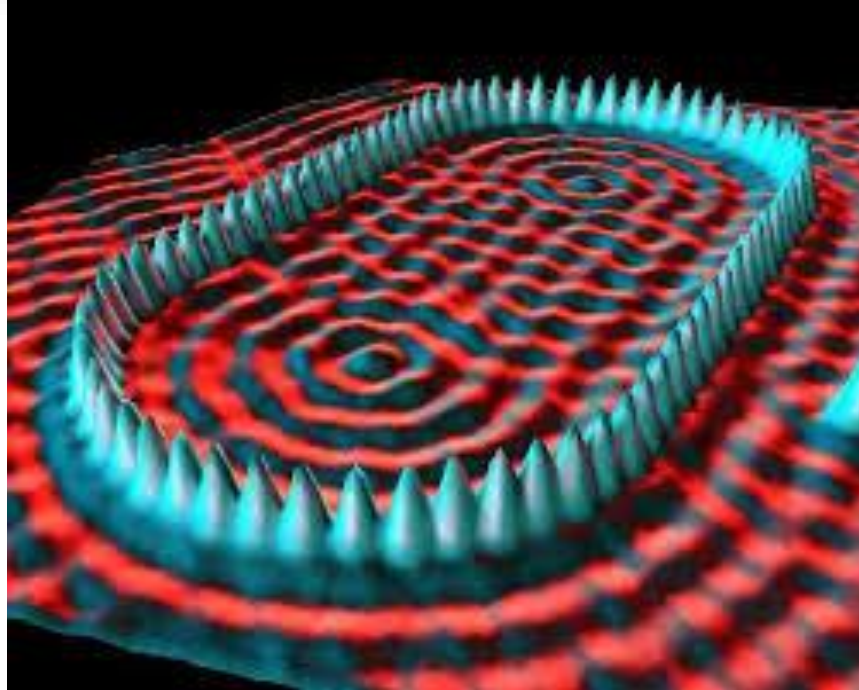
Standing waves: a way to “see” and measure electrons

Ag atoms on Ag(111)



R. Berndt, Uni Kiel

Standing waves: a way to “see” and measure electrons



Scattering coefficients

Heller Nature **369**, 464 (1994)

Quantum mirage

Manoharan Nature **403**, 512 (2000)

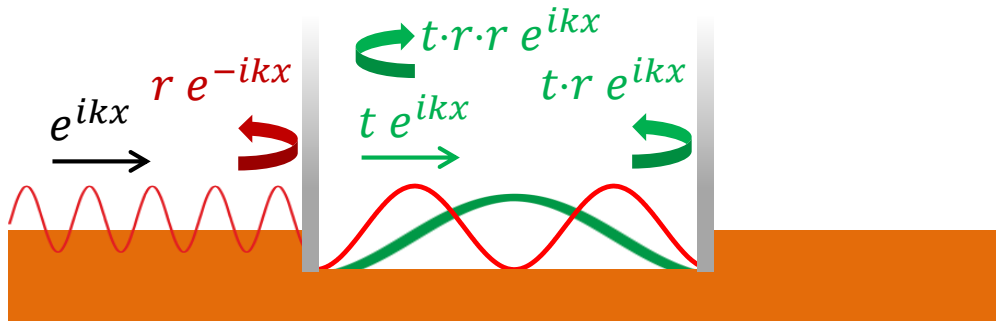
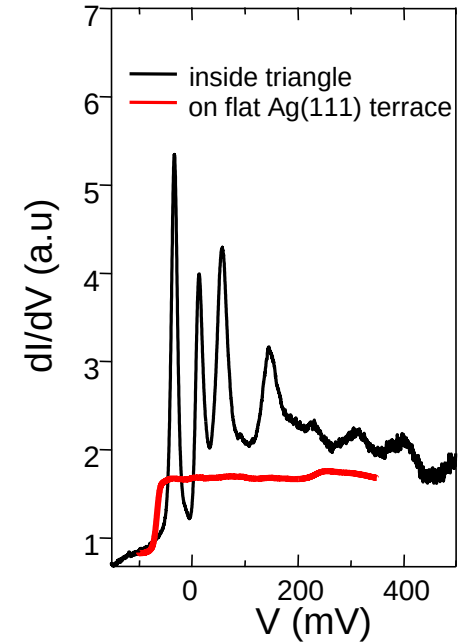
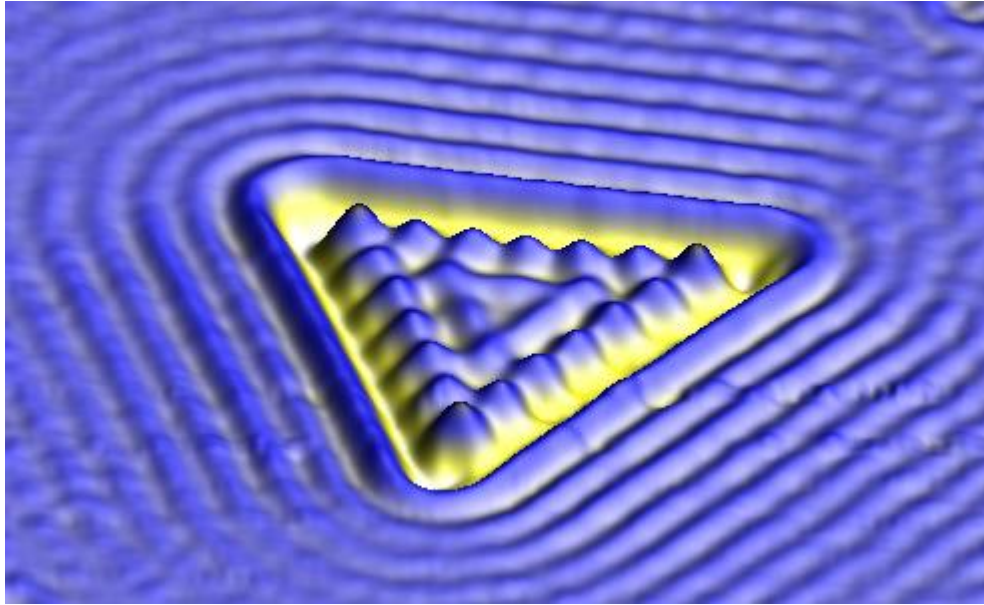
Quantum invisibility

Fransson Nano Lett. **10**, 1600 (2010).

Standing waves: a way to “see” and measure electrons

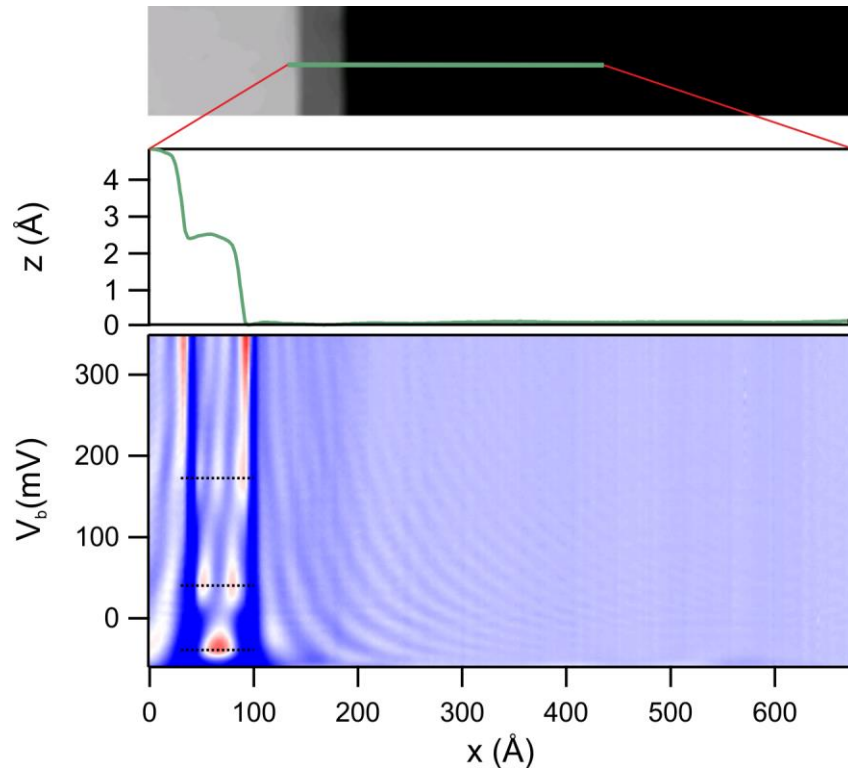
Vacancy island on Ag(111)

Top



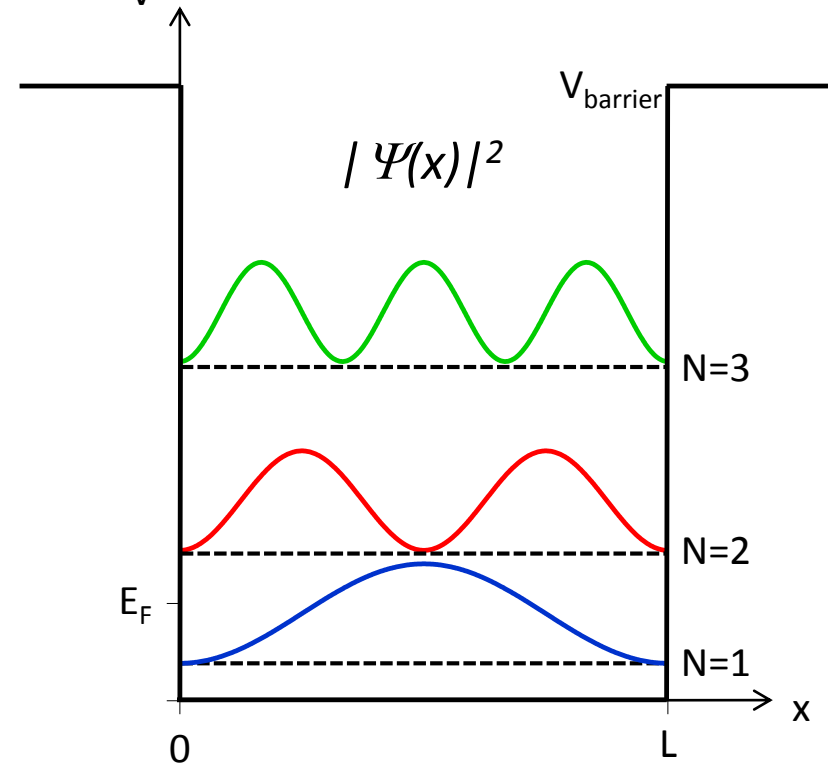
Electron scattering and confinement in 1D: monoatomic steps

Monoatomic steps on Ag(111)



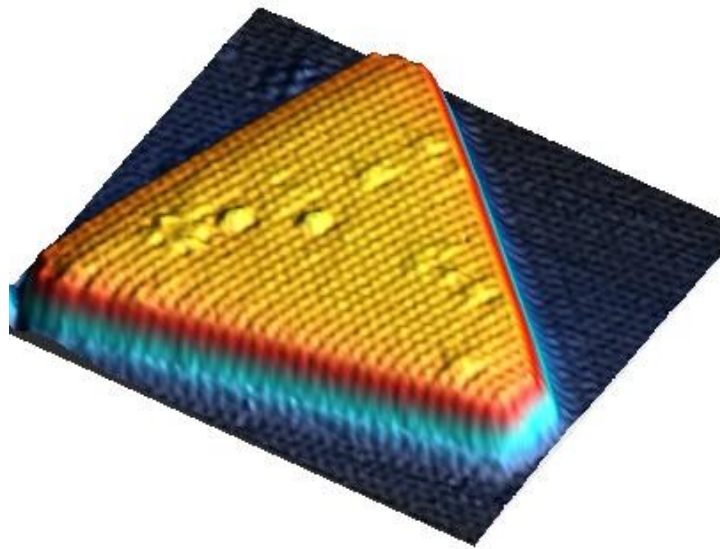
A. Mugarza et al. PRL **87**, 3 (2001)
F. Schiller et al., NJP **16**, 123025 (2014)
J. E. Ortega et al. PRB **87**, 115425 (2013)
J. E. Ortega et al., Handbook of Surface Science
(Springer, coming son!)

Particle in a box

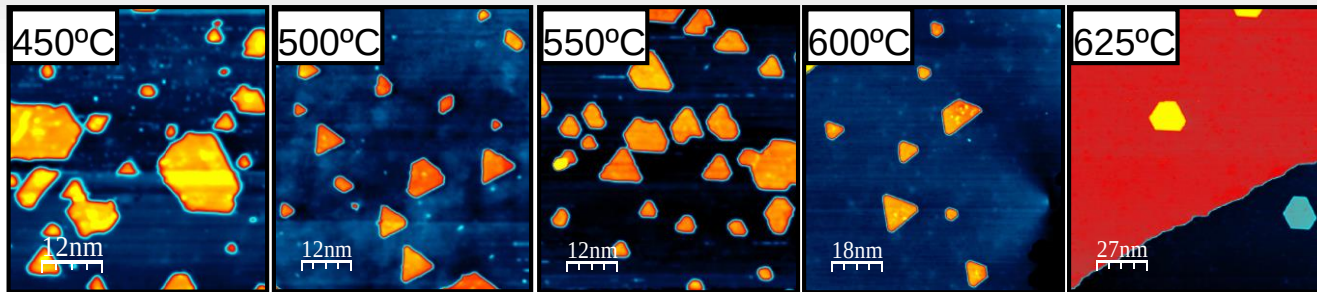


Quantum interference effects in a potential well: standing waves and discrete energy levels

Graphene nanoislands

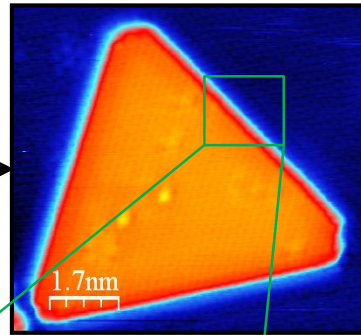


Synthesis of nanographene on Ni(111)

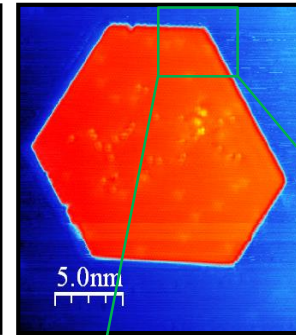


Propene dose at RT
 Anneal to: 450°C irregular clusters
 500°C triangles
 625°C hexagons

nanotriangles



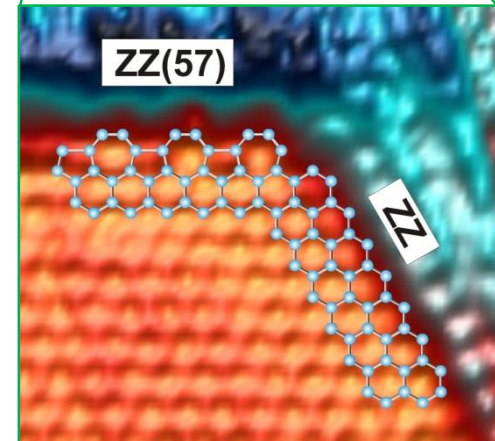
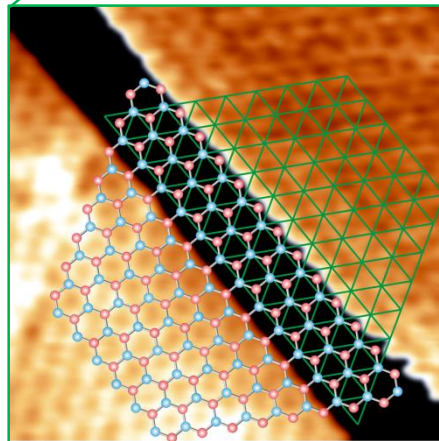
nanohexagons



M. Ollé Nano Lett. 12 (2012)

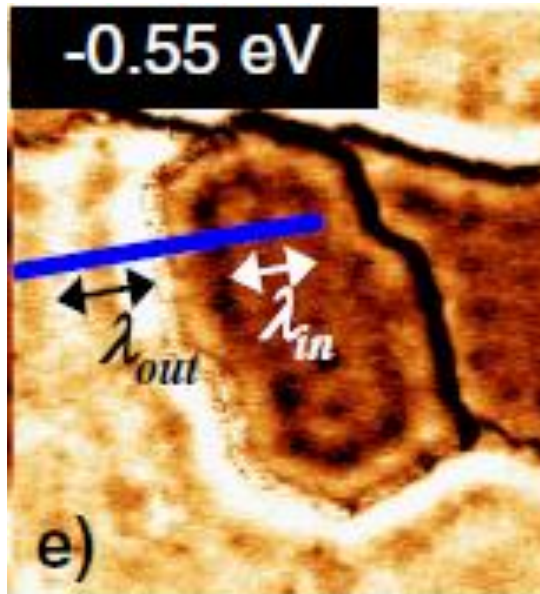
- ☐ Stacking: top-fcc
- ☐ Triangles: zig-zag hollow
- ☐ Hexagons: edges with different reconstruction

*A. Garcia-Lekue et al.
 JPCC 119, 2015*

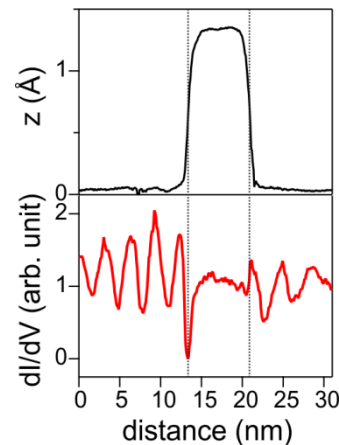


Surface state scattering at graphene

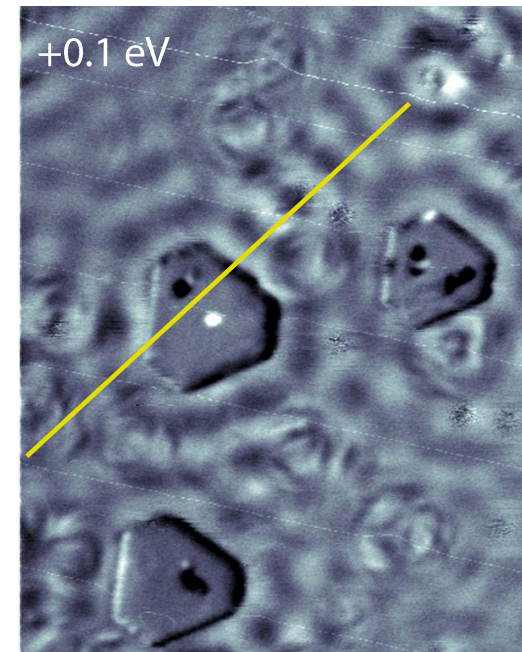
G/Ir(111)



S. J. Altenburg et al., PRL 108 (2012)
D. Subramaniam et al., PRL 108 (2012)
S. K. Hämäläinen et al., PRL 107 (2011)

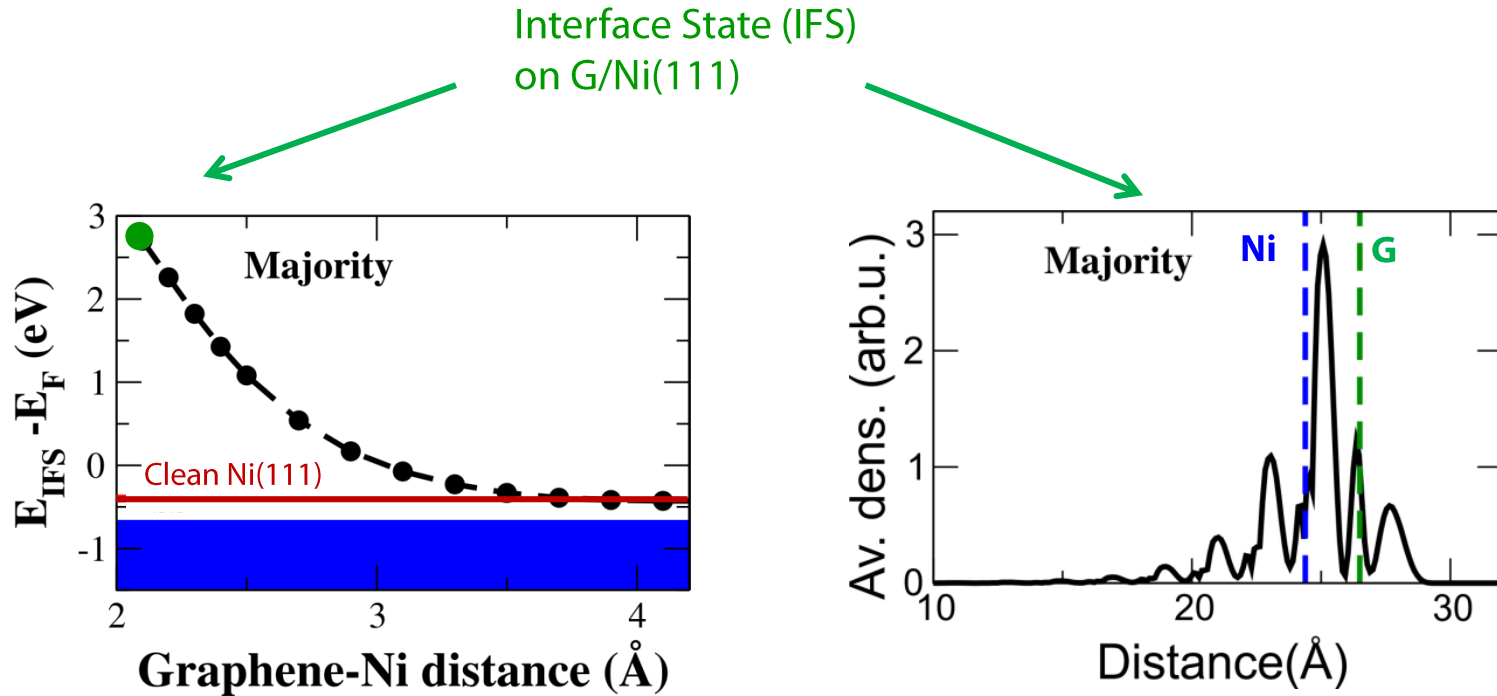


G/Ni(111)

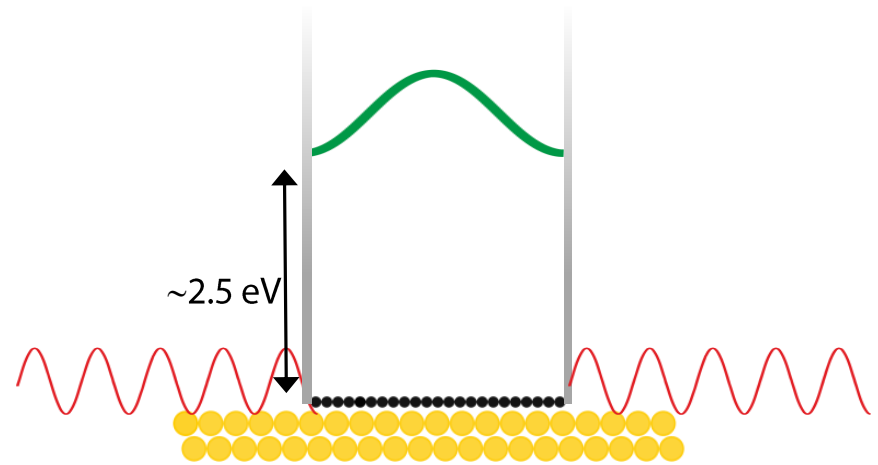
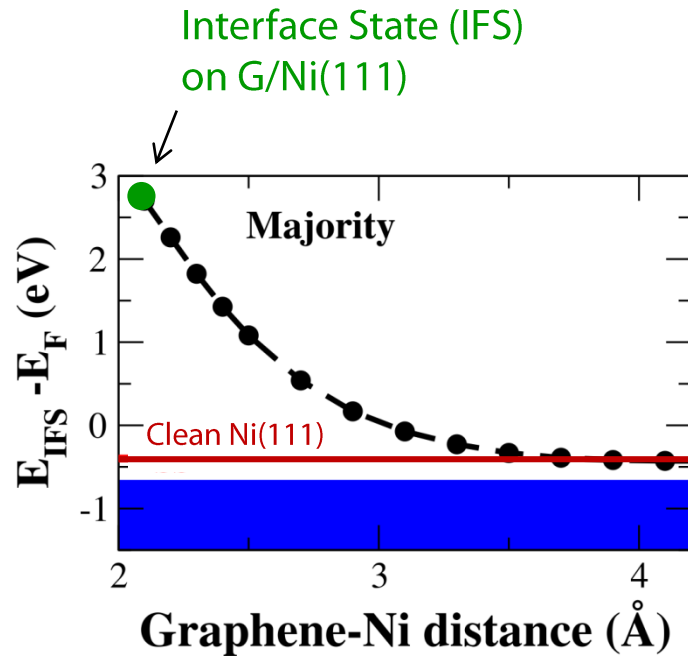


A. Garcia-Lekue et al., PRL 112, 066802 (2014)

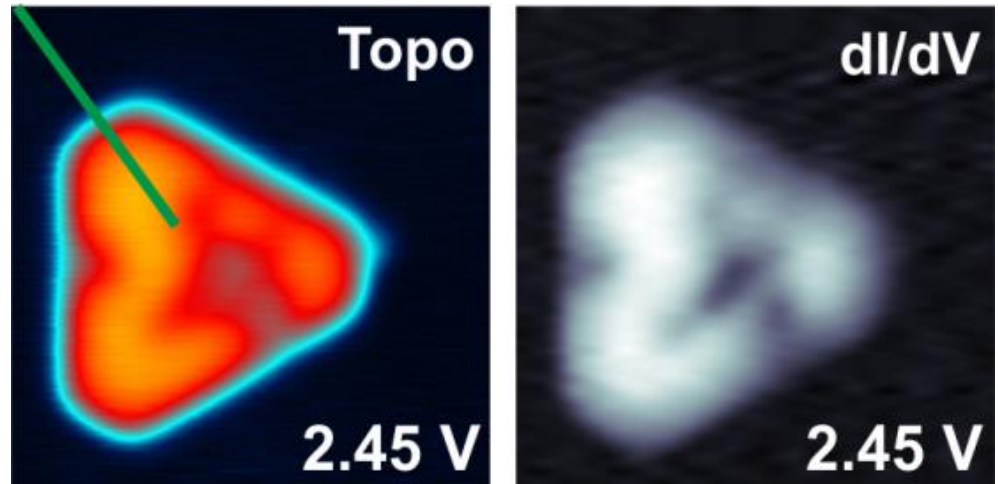
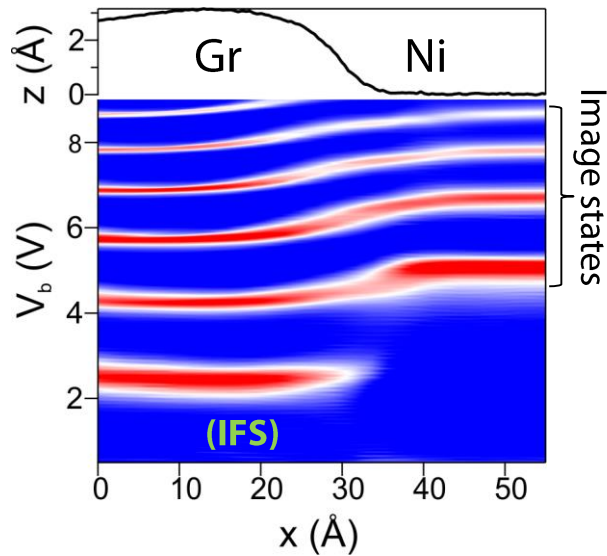
Surface state scattering at graphene



Surface state scattering at graphene

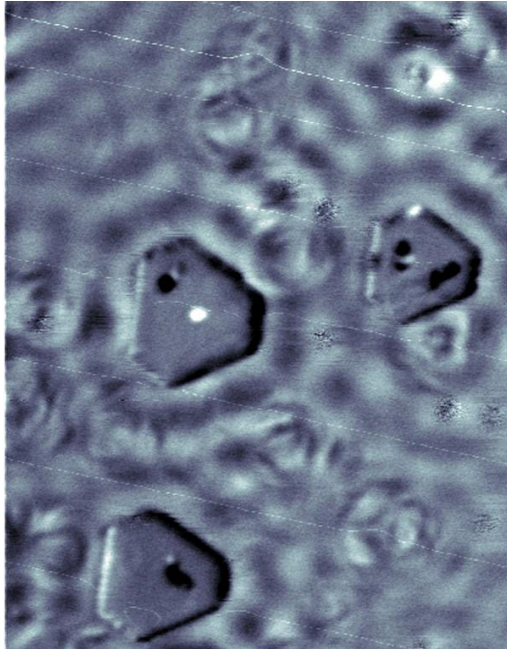


Surface state scattering at graphene

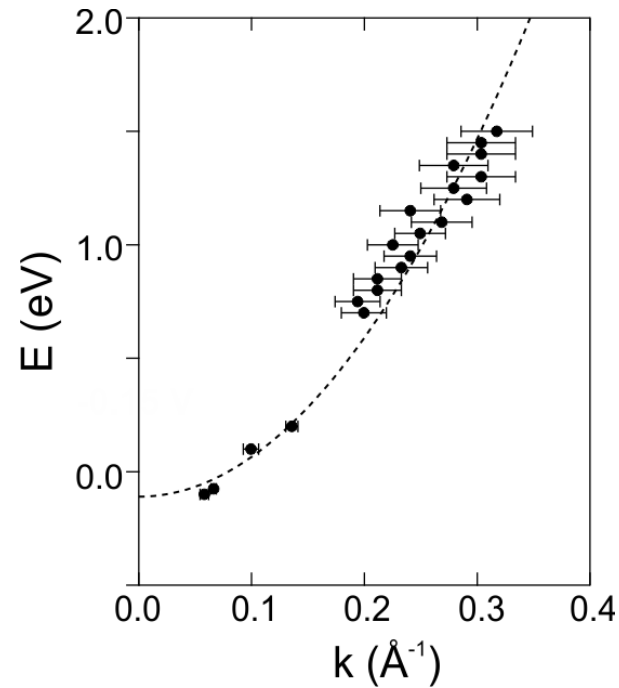


Lateral spin filter

dI/dV map

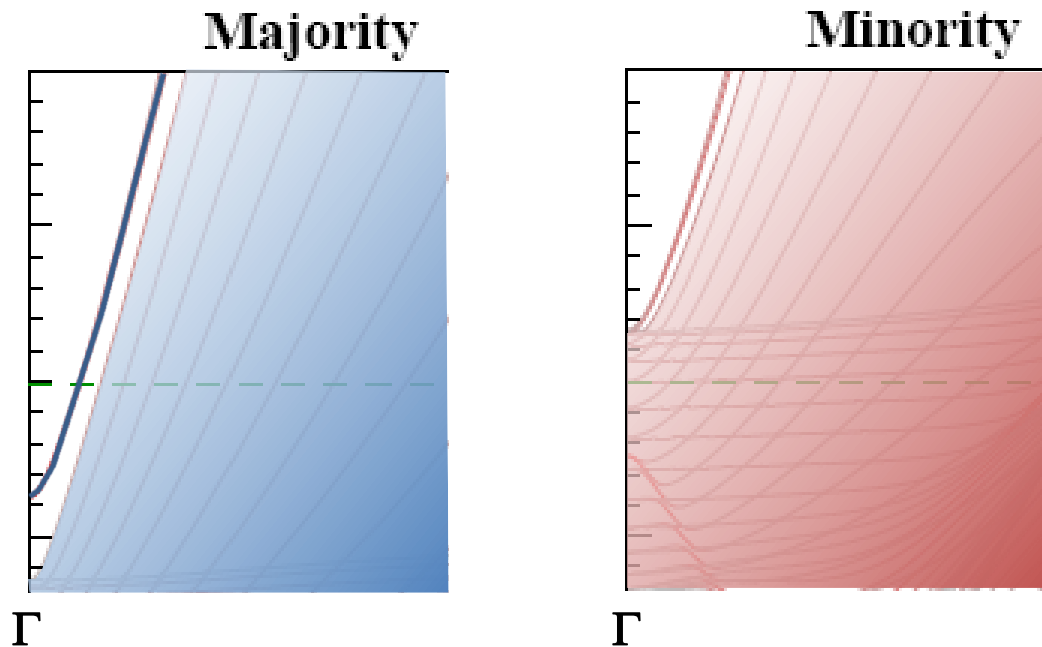


Dispersion obtained from periodicity of standing waves

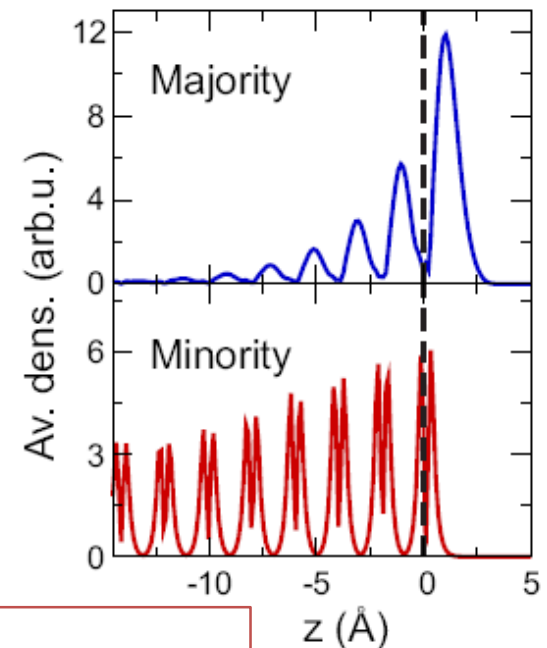


Graphene: perfect reflector for Ni surface states
...for both spins?

Coupling to bulk states



Planar average of density of states of S_1 at Γ



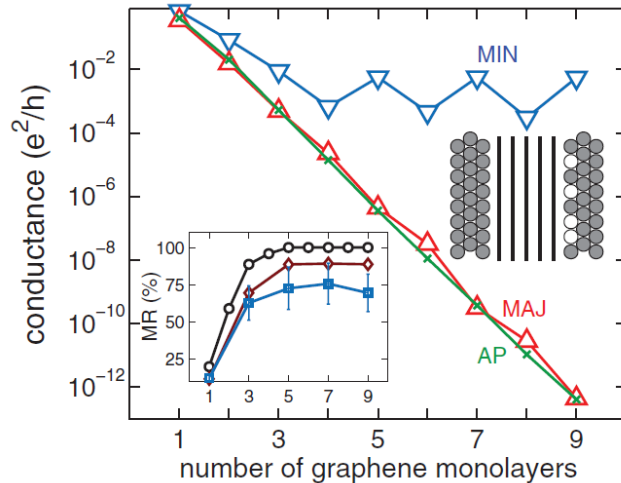
Minority band:

- ❑ Closer to bulk states $\rightarrow$ larger absorption to bulk
- ❑ Larger bulk penetration $\rightarrow$ less sensitivity to surface potentials



Spin-dependent scattering

Coupling to bulk states

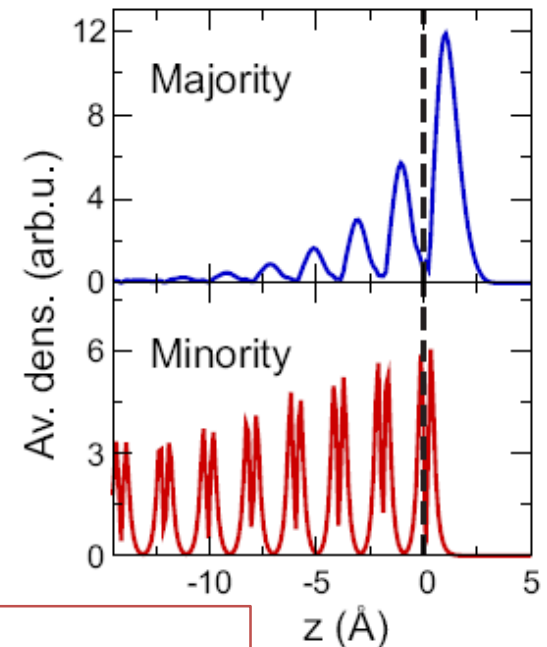


Spin filtering in G/Ni interfaces:

- Magnetoresistance close to 100%

V. M. Karpan et al. PRL 99 (2007)

Planar average of density of states of S_1 at Γ



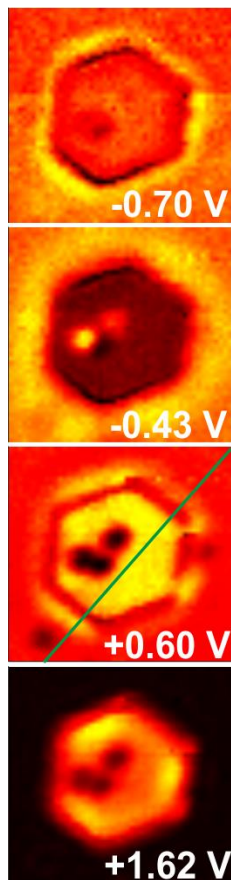
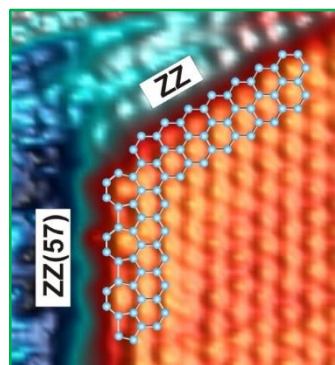
Minority band:

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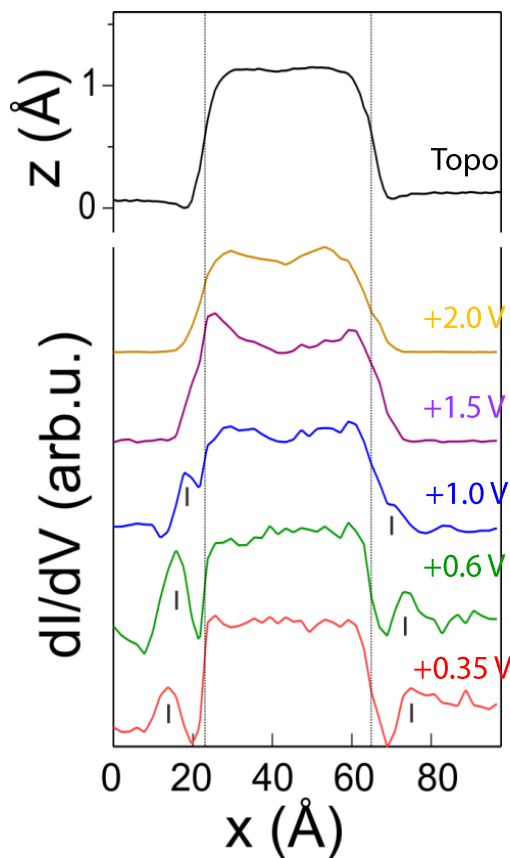
↓
Spin-dependent scattering

Effect of graphene edge on scattering

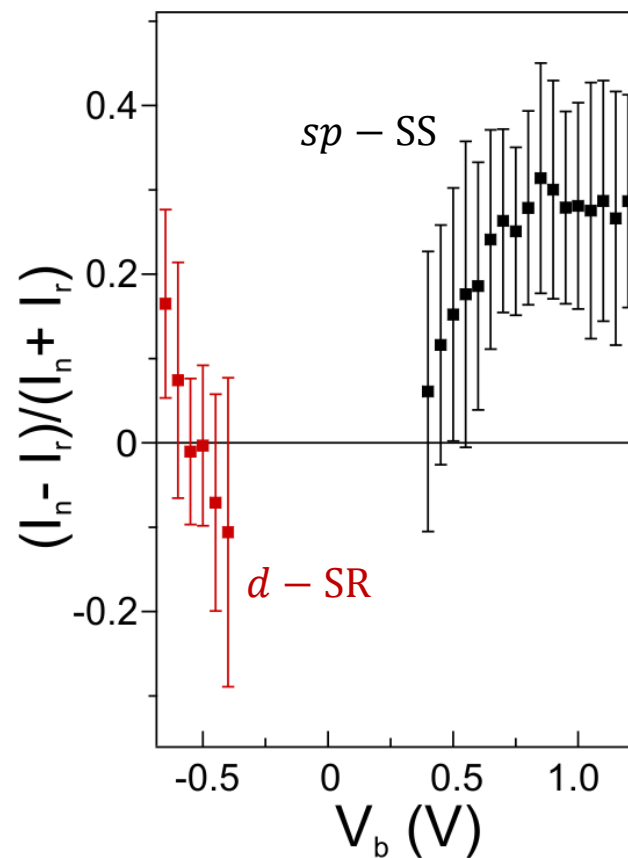
dl/dV maps



dl/dV profiles at different energy

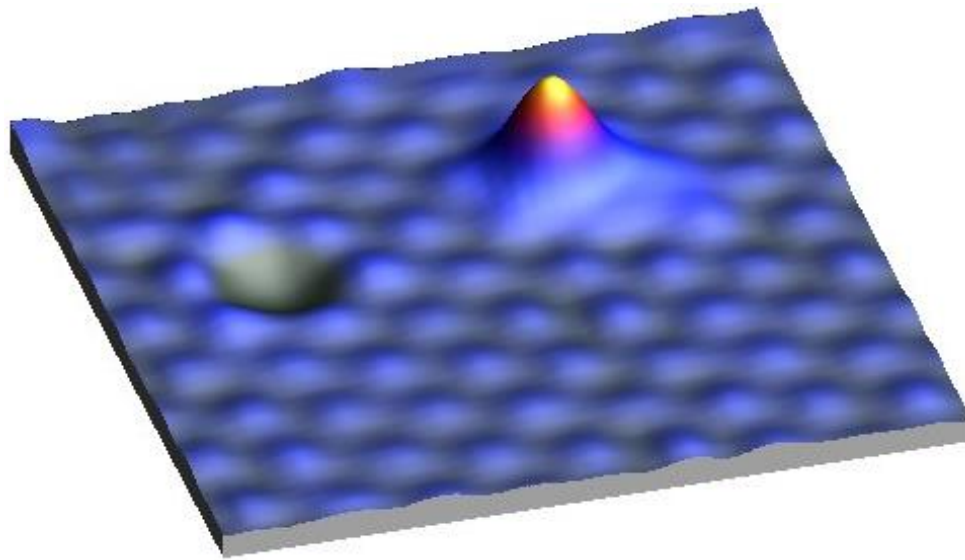


Asymmetry of scattering amplitude

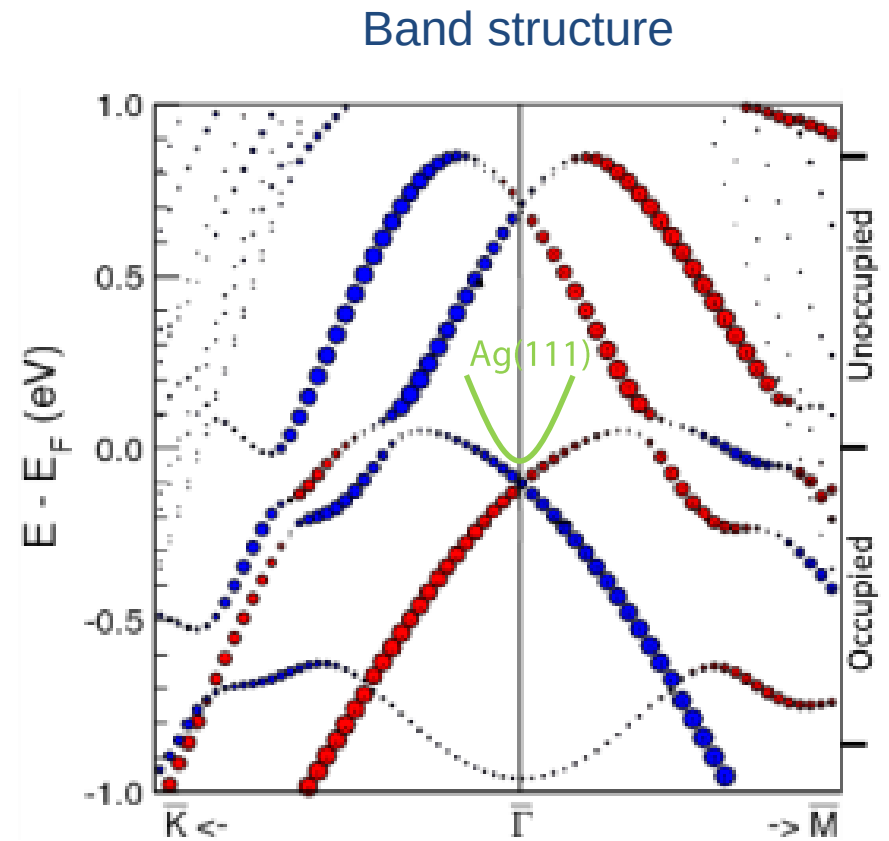
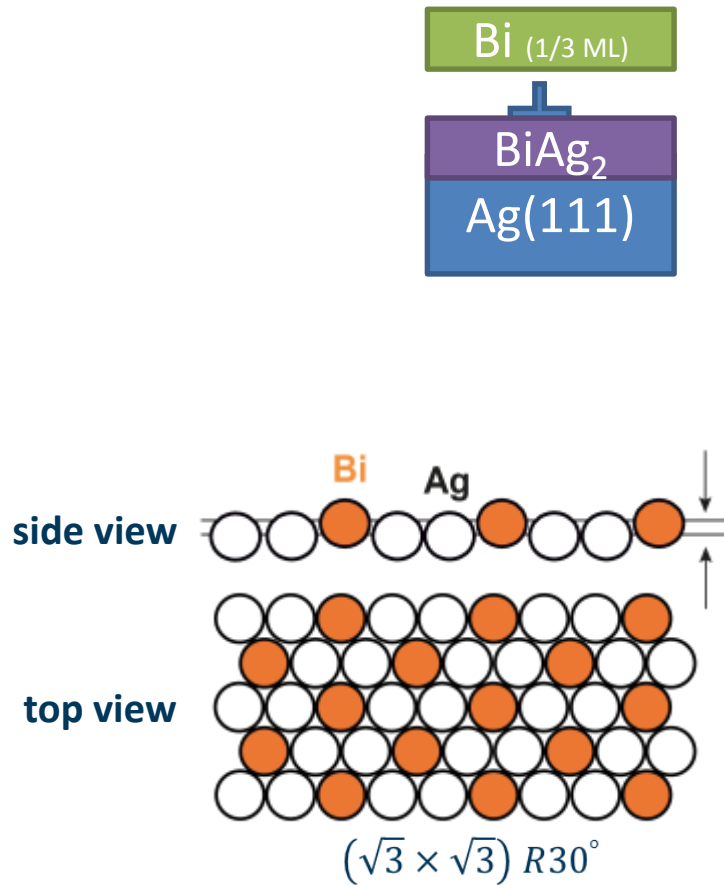


- Scattering stronger in zig-zag edges
- Asymmetry depends on energy and type of surface state

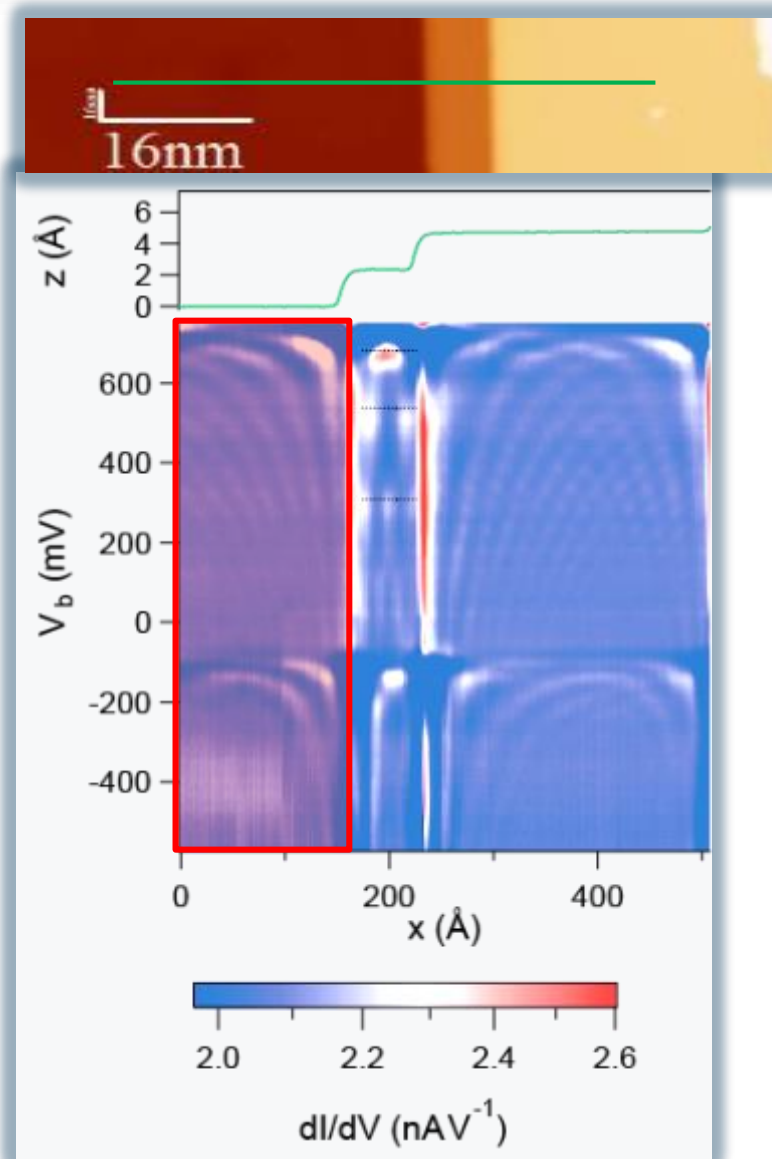
Heavy metal alloys



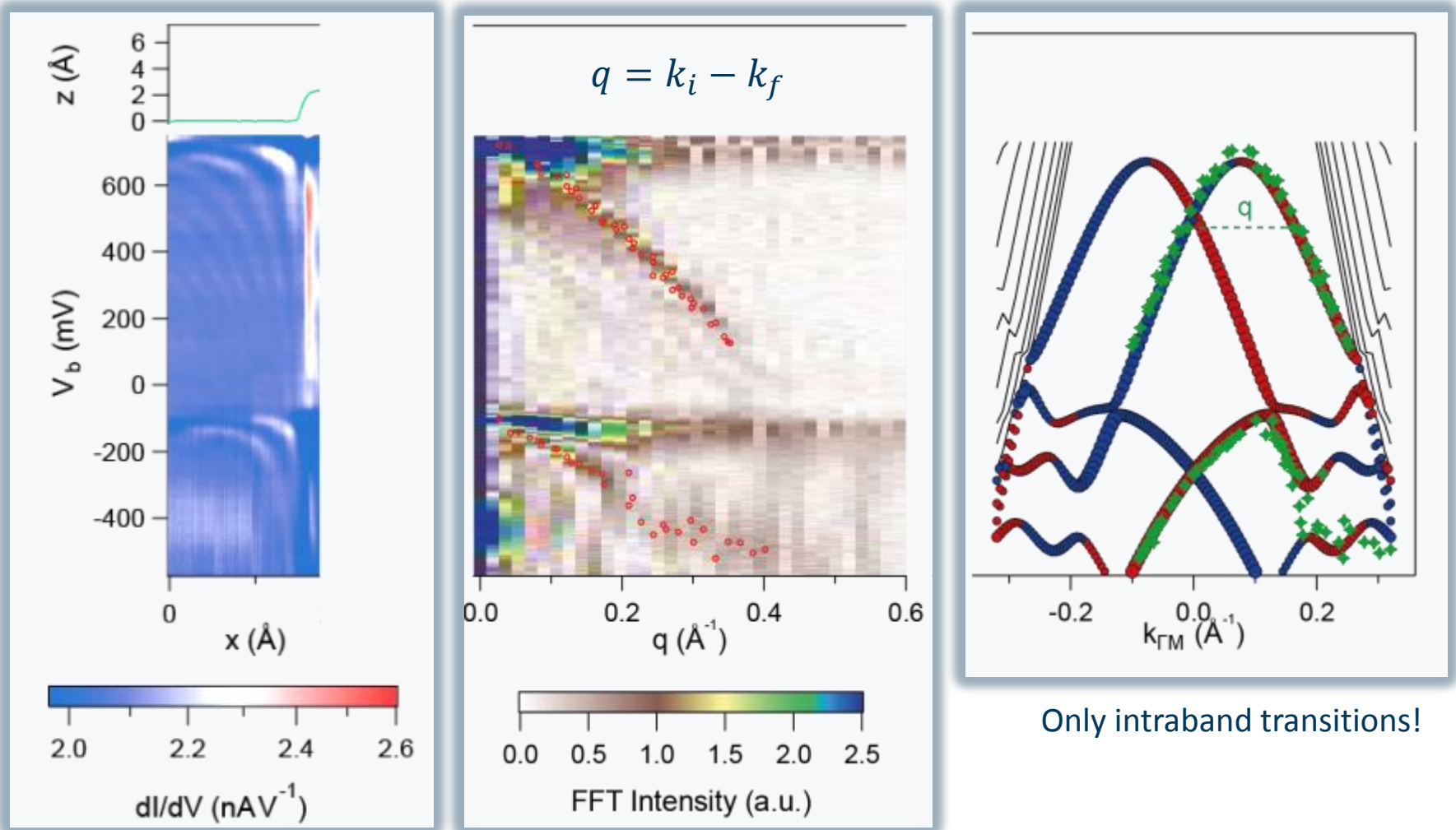
BiAg₂ surface alloy



Scattering at semi-infinite terraces

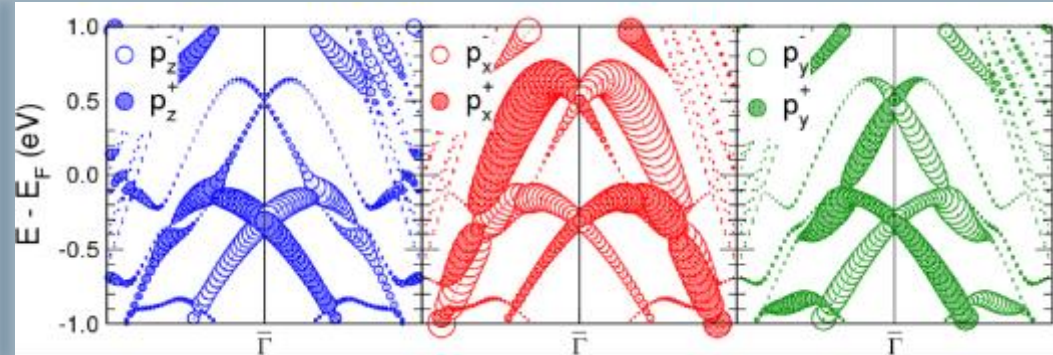
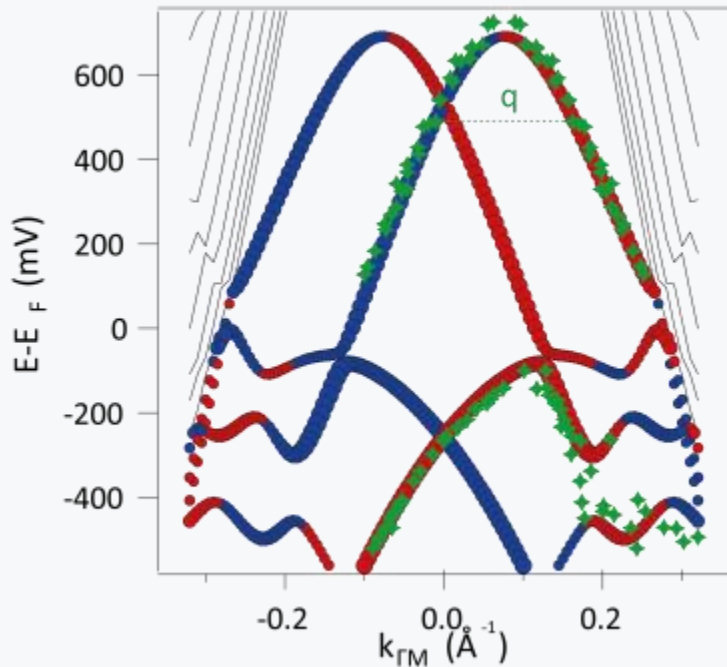


Scattering at semi-infinite terraces



Only intraband transitions!

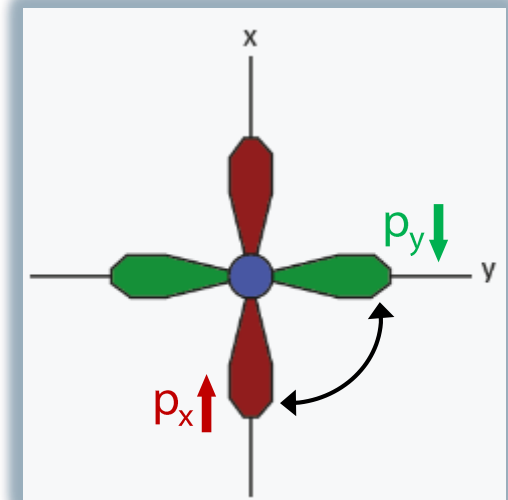
Spin-orbital flip scattering



p_z

p_x

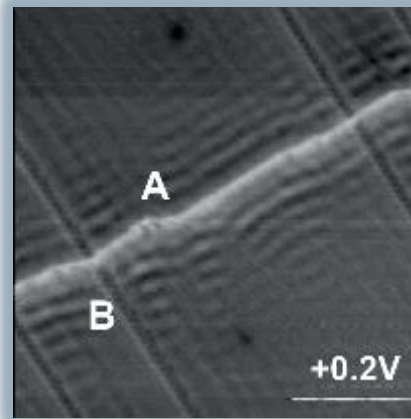
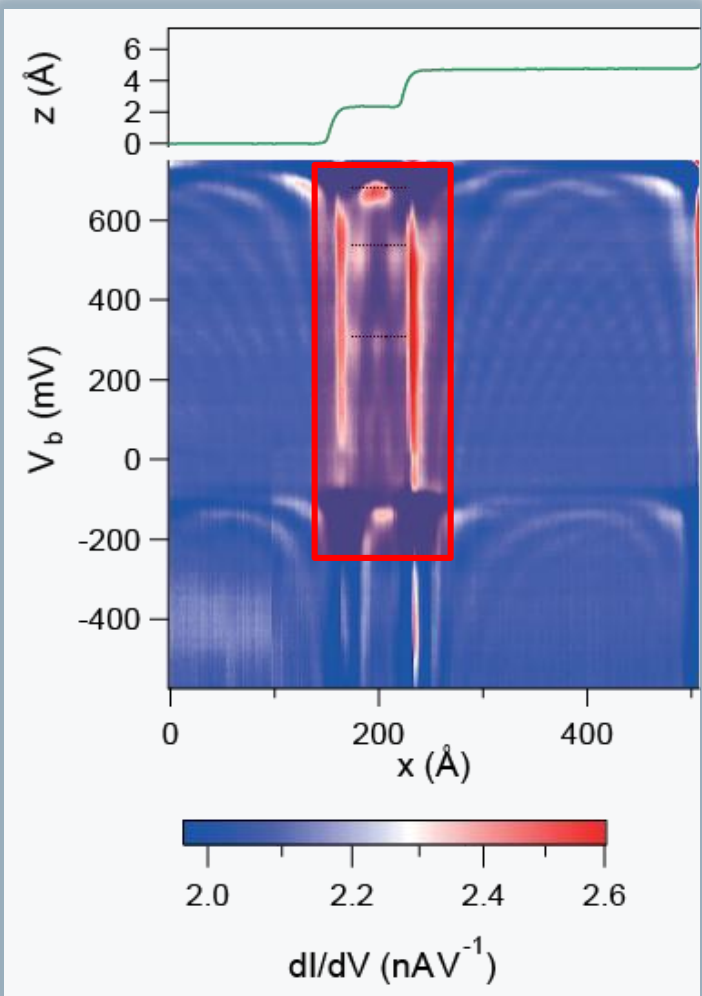
p_y



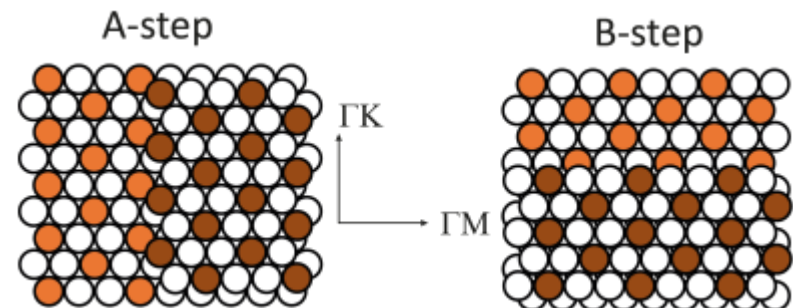
$$\langle p_{y,\downarrow} | \mathbf{L} \cdot \mathbf{S} | p_{x,\uparrow} \rangle \neq 0$$

$$\Delta \mathbf{J} = \Delta \mathbf{L} + \frac{1}{2} \Delta \mathbf{S} = 0$$

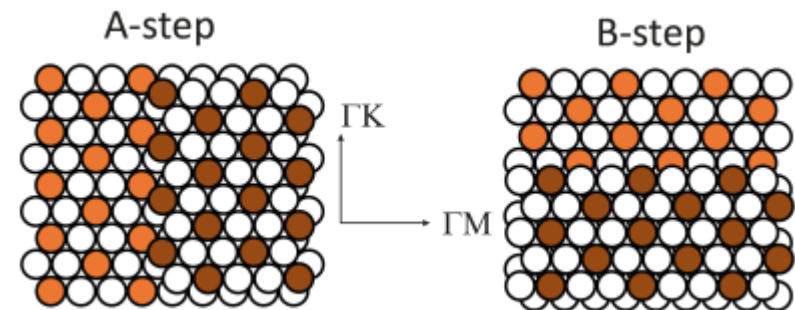
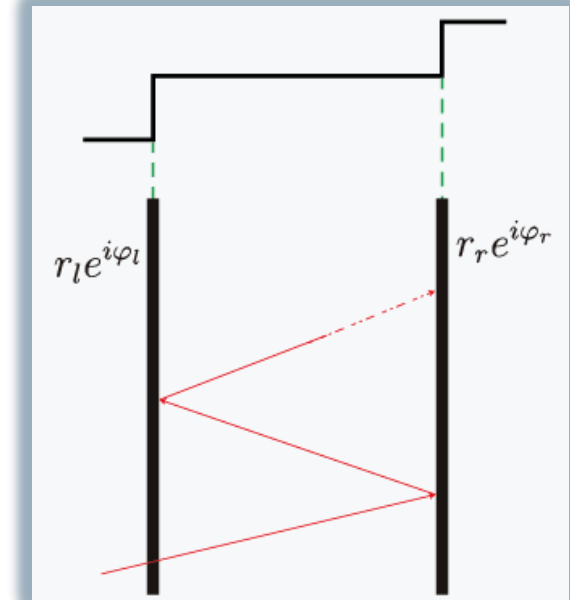
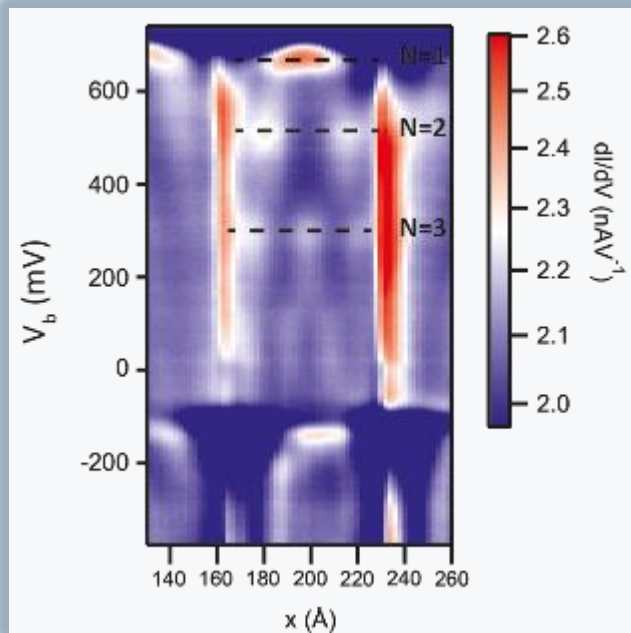
Scattering in resonators



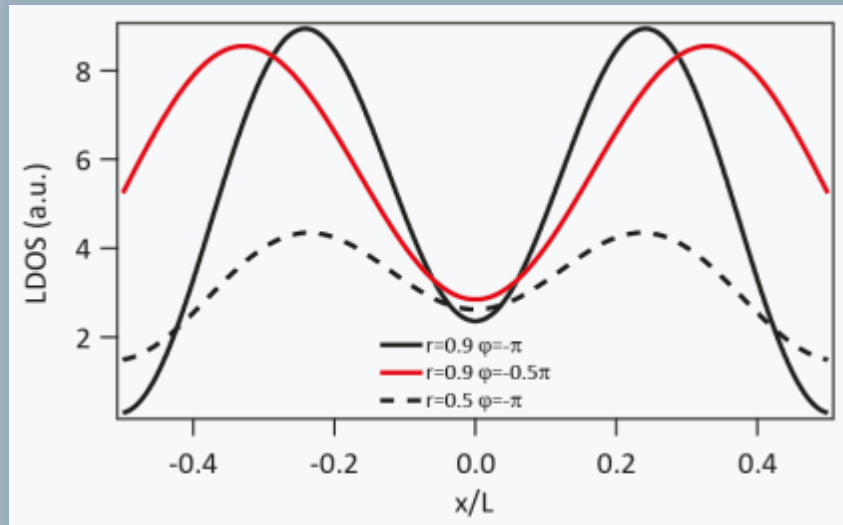
Different intensity



Scattering in resonators

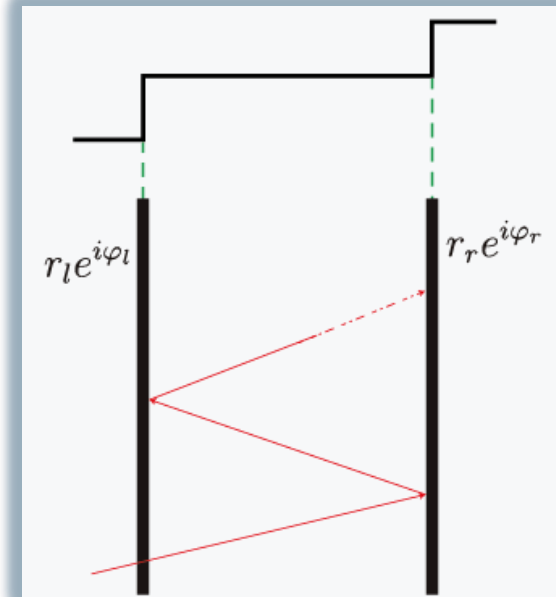


Fabry Perot resonator



Infinite barrier $\rightarrow \varphi = -\pi$

L. Bürgi et al., Phys. Rev. Lett. 81, 5370 (1998).



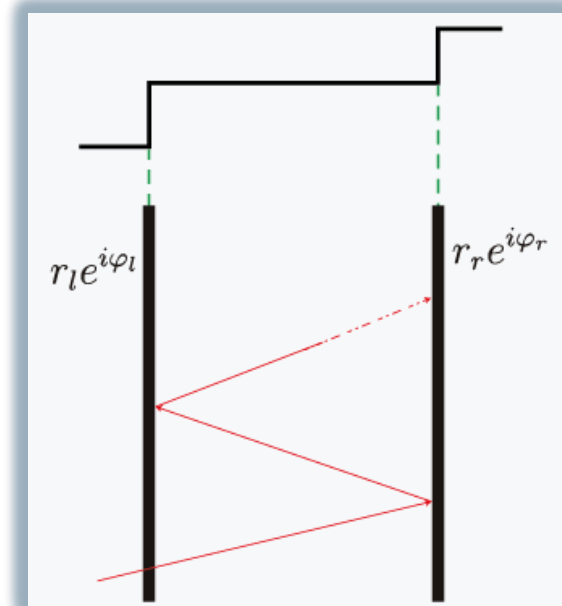
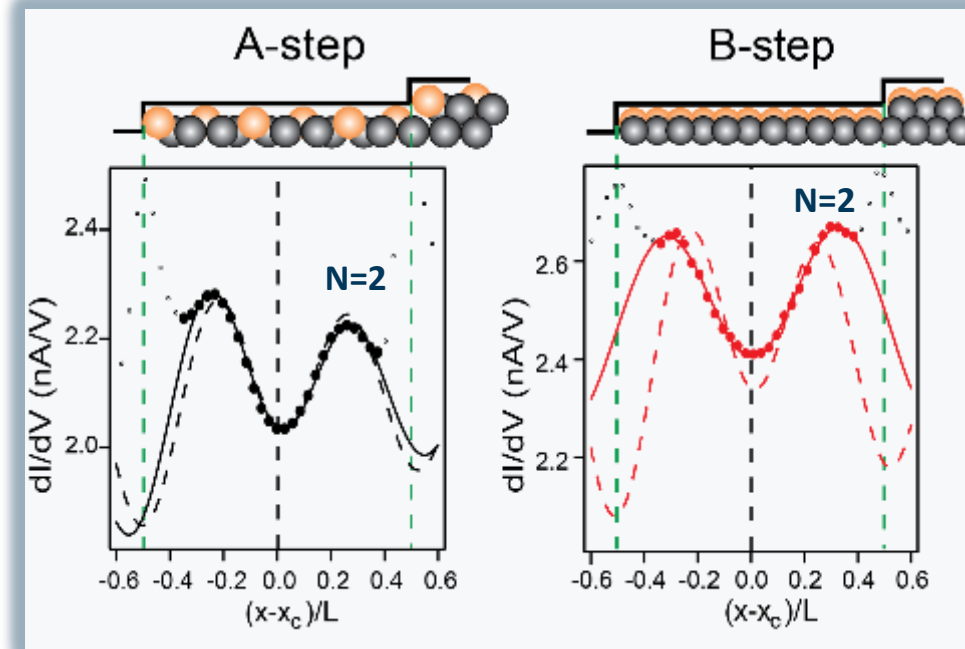
$$\frac{dI}{dV}(E, x) \propto \rho_b + \frac{\rho_s}{\pi} \int_{k_R}^{k_E} dk_x \frac{1}{\sqrt{k_E^2 - (q/2)^2}} \times \left\{ \frac{(1 - r_l^2) [1 + r_r^2 + 2r_r \cos(qL(x - \frac{1}{2}) - \varphi_r)]}{1 + r_l^2 r_r^2 - 2r_r r_l \cos(qL + \varphi_r + \varphi_l)} + \frac{(1 - r_r^2) [1 + r_l^2 + 2r_l \cos(qL(x + \frac{1}{2}) + \varphi_l)]}{1 + r_l^2 r_r^2 - 2r_r r_l \cos(qL + \varphi_r + \varphi_l)} \right\}$$

$$k_E = k - k_R = \sqrt{2m_{\text{eff}}/\hbar^2(E - E_0)}$$

$$q = 2(k_x - k_R)$$

S. Schirone et al. PRL 114, 166801 (2015)

Effect of resonator edge structure

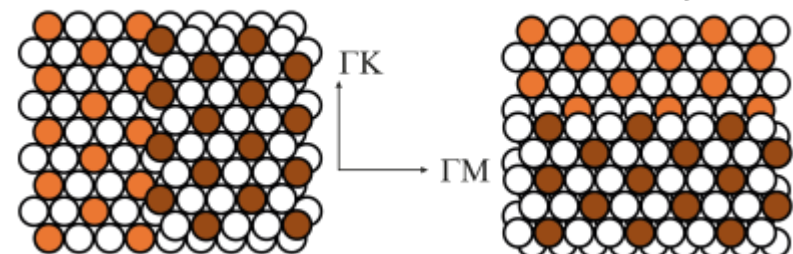


strong resonator

weak resonator

A-step

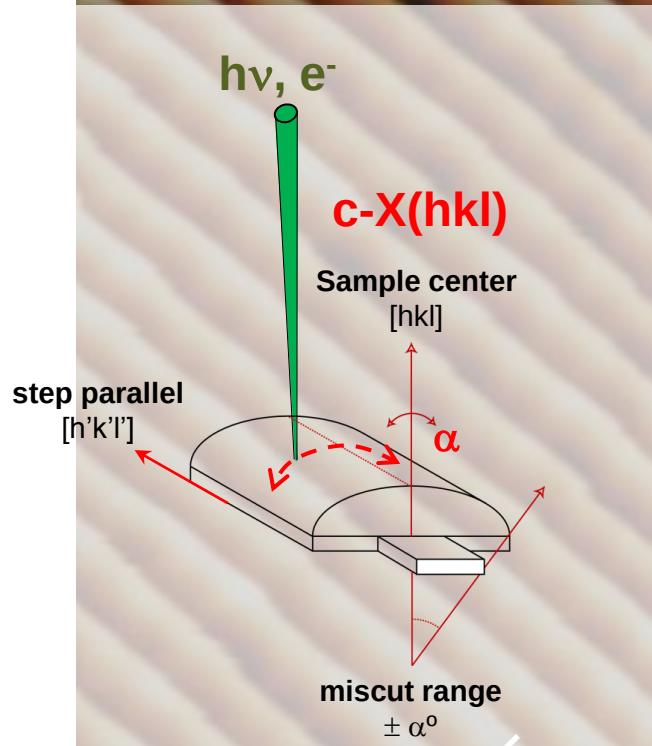
B-step



	r_l	r_r	φ_l/π	φ_r/π
A-step	0.70 ± 0.30	0.21 ± 0.07	-0.83 ± 0.02	-0.95 ± 0.02
B-step	0.40 ± 0.18	0.36 ± 0.15	-0.56 ± 0.02	-0.51 ± 0.02

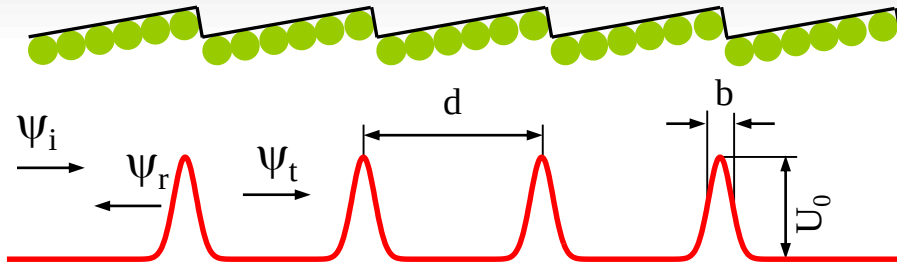
The Curved Surface Approach

Curved surfaces allows to **tune** and achieve **atomic scale control** over **physical-chemical** processes related to **atomic steps**



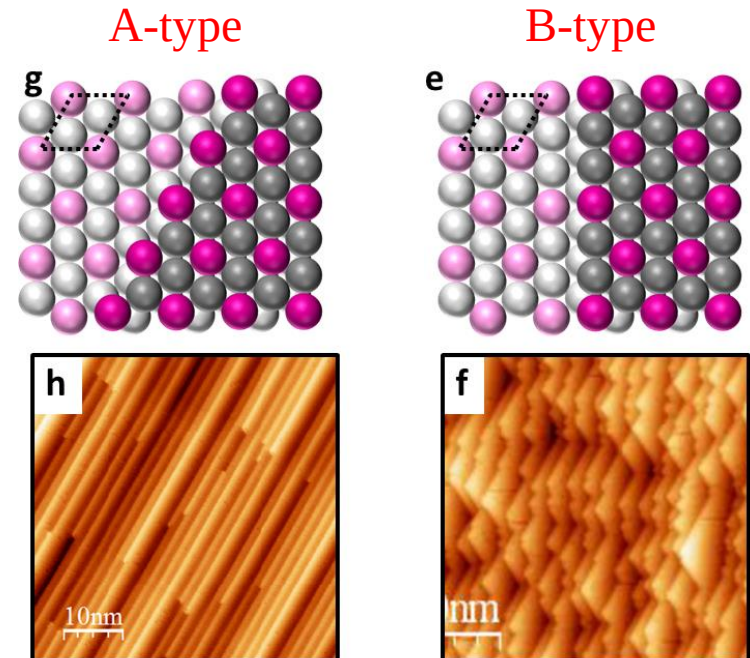
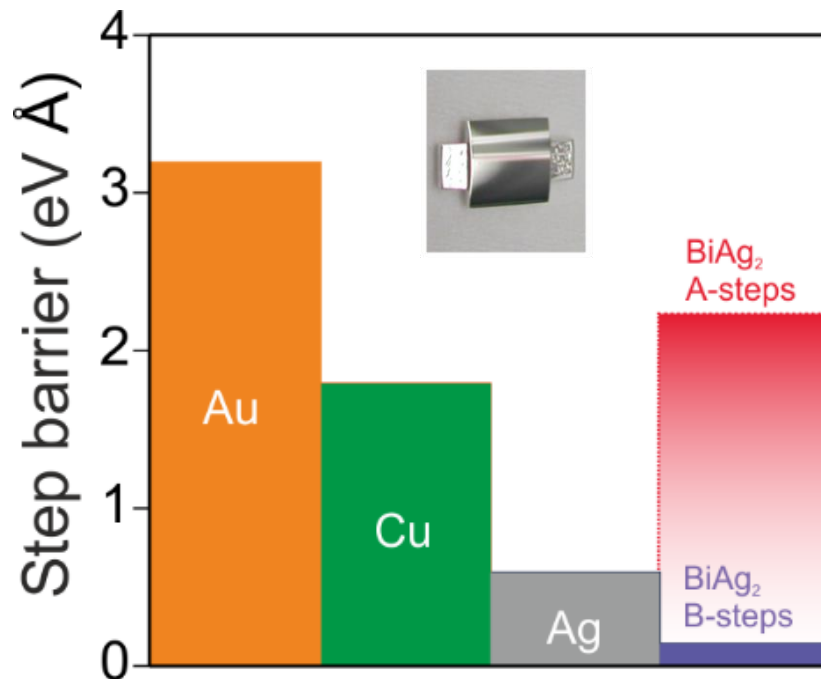
d
Tunable

Scattering analysis framework: 1D Kronig-Penney model



$$\Delta E = \frac{2\hbar^2}{m^*} \frac{1}{d^2} \left(\text{tg}^{-1} \left(\frac{U_0 b}{\sqrt{\frac{2\hbar^2}{m^*} \Delta E}} \right) \right)^2$$

1D Kronig-Penney model



Atomic Manipulation and Spectroscopy Group

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D. Sanchez-Portal
J. E. Ortega
F. Schiller
E. Krasovskii

CFM
CFM
CFM
CFM

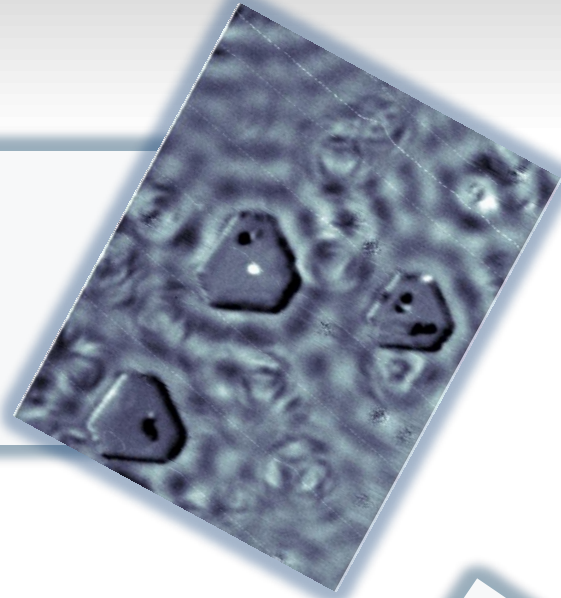
dipco

J. Lobo-Checa
 **icma**
Instituto de Ciencia
de Materiales de Aragón

P. Gambardella
ETH zürich

G. Bihlmayer
 **JÜLICH**
FORSCHUNGSZENTRUM

- Lateral spin-filtering
- Efficiency depends on edge structure



- New scattering mechanism in spin-orbitronic materials
- Controlled by chemical composition of steps

